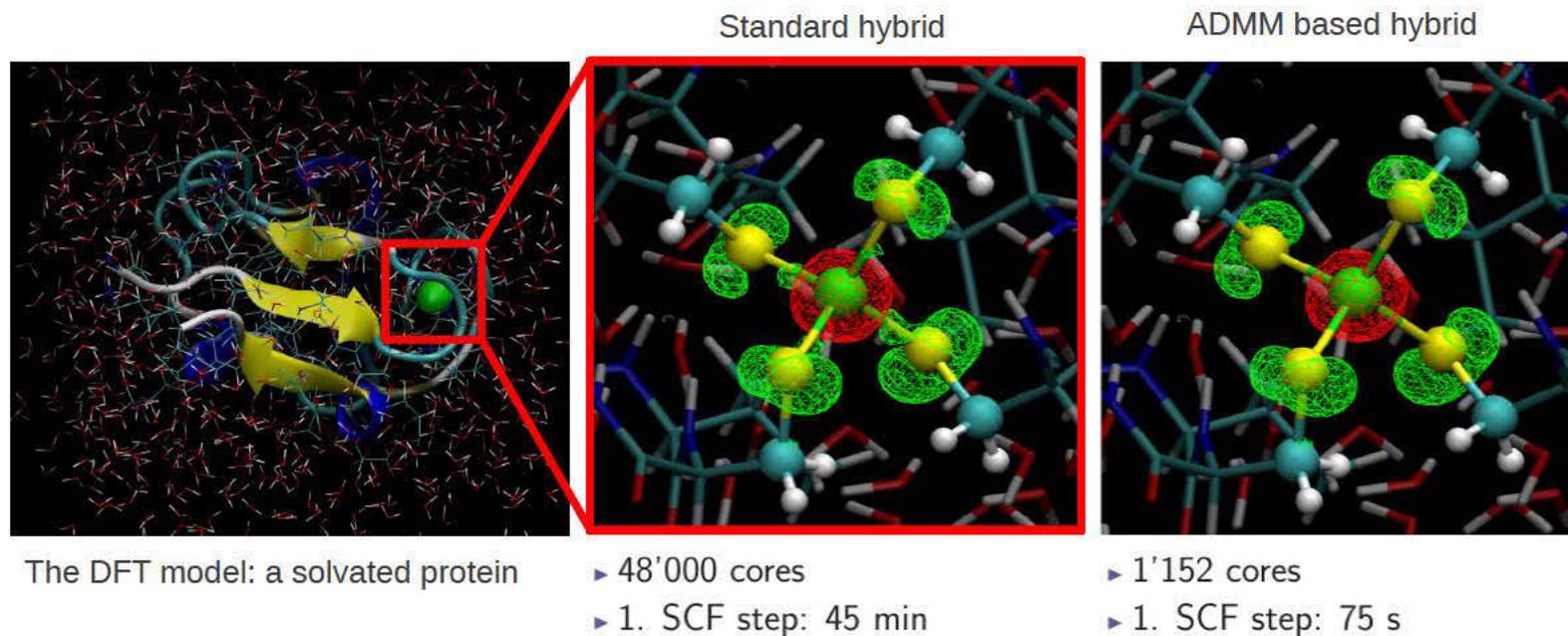


# Large scale DFT based Molecular Dynamics in the condensed phase

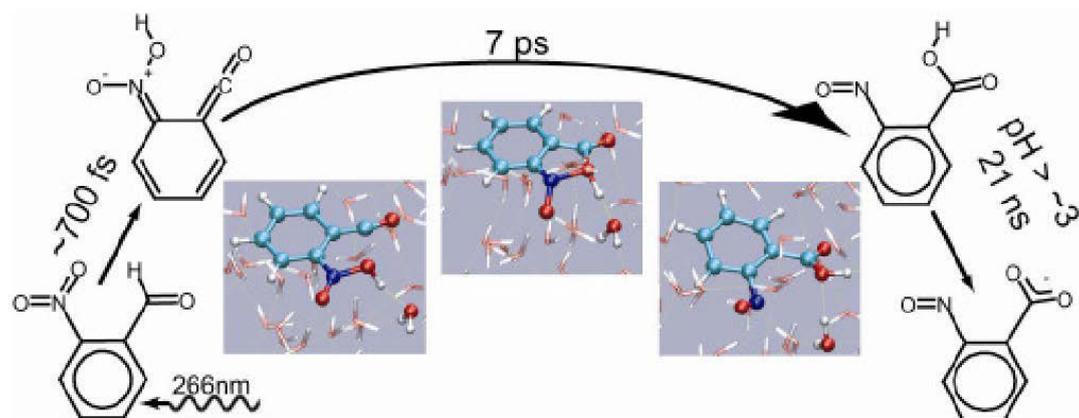


Joost VandeVondele  
Nanoscale Simulations, ETH Zurich

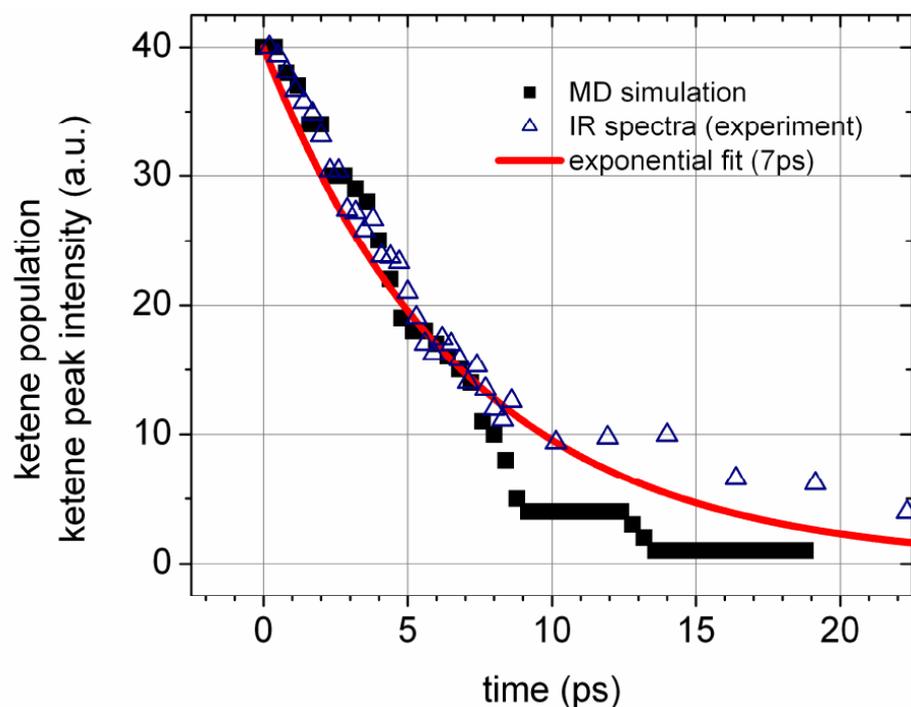
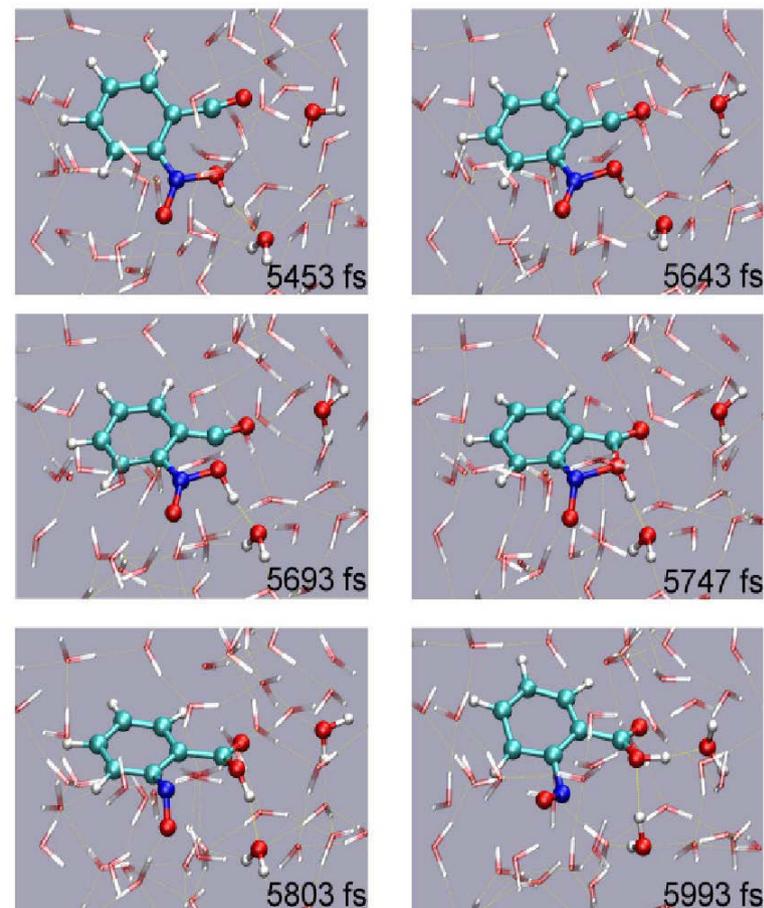
# Ab Initio Molecular Dynamics

- Molecular Dynamics
  - Classical equations of motion for the ions:  $F=ma$
  - NVE and other ensembles (NVT, NPT)
  - Crucial for systems including liquids
- 'Ab Initio' or 'First Principles'
  - Forces on atoms derive from a non-empirical description of the electronic structure
  - 99% DFT

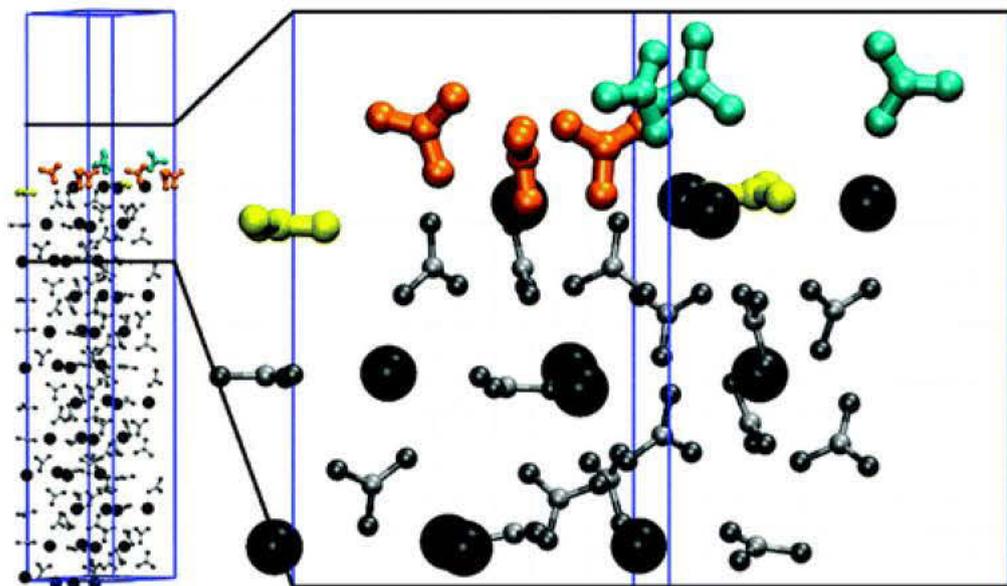
# An chemist's example: oNBA



ortho-Nitrobenzaldehyde

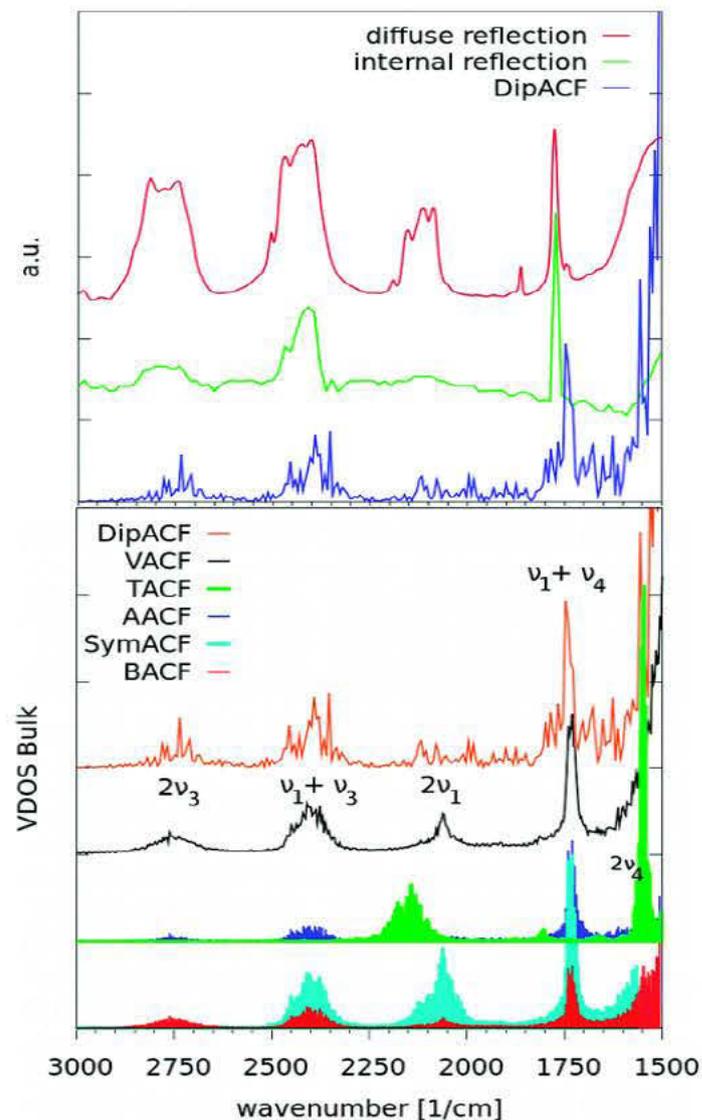


# Solids: spectroscopy at interfaces

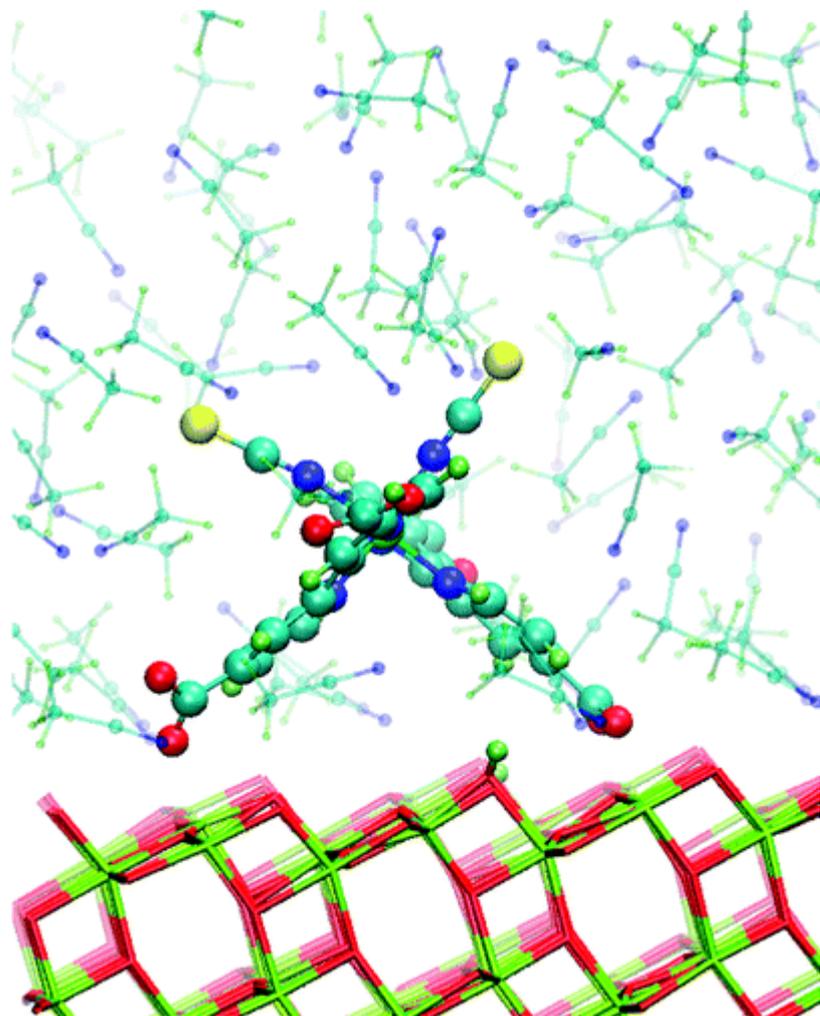


BariumNitrate ( $\text{Ba}(\text{NO}_3)_2$ ) :  
a catalyst for automotive applications ( $\text{NO}_x$ )

Ab initio MD can be used to capture non-harmonic effects  
Such as overtones and combination bands.



# Interfaces at ambient conditions



We want to build and study models that realistically include solids, liquids, molecules

An example interface, as found in Dye sensitized solar cells.

# AIMD is GGA dominated

In order for AIMD to be practical,  
the SCF + forces (BOMD) need to be computed  
in O(60sec) for typical systems O(100 atoms)

20000 MD steps of 0.5 fs \* 60sec / step = 2 weeks for 10 ps

- Highly efficient schemes for GGAs are needed
- Highly robust methods are essential

# AIMD is for short timescales

10ps – 100ps is the typical lengthscale of AIMD

- Only very fast processes can be observed spontaneously
- Activated events (reactions) or slow dynamics (diffusion) need more advanced techniques
  - Free energy methods (constraints, enhanced sampling, etc... needed).

Efficient GGA DFT

CP2K

# Gaussian and plane waves: GPW in CP2K

- **Primary basis: Gaussians** → Chemistry
  - compact
  - sparse  $H^{ks}$  (and  $P$ )
  - Many terms analytic
- **Auxiliary basis: Plane waves** → Physics
  - regular grid for  $e^-$  density
  - FFT for Poisson equation
  - No four center integrals needed

The GPW algorithm : compute the GGA Kohn-Sham matrix  
in  $O(N)$  time, PBC are natural.

# Basic Computational local DFT

$$n(r) = \sum_{\mu\nu} P^{\mu\nu} \varphi_{\mu}(r) \varphi_{\nu}(r)$$

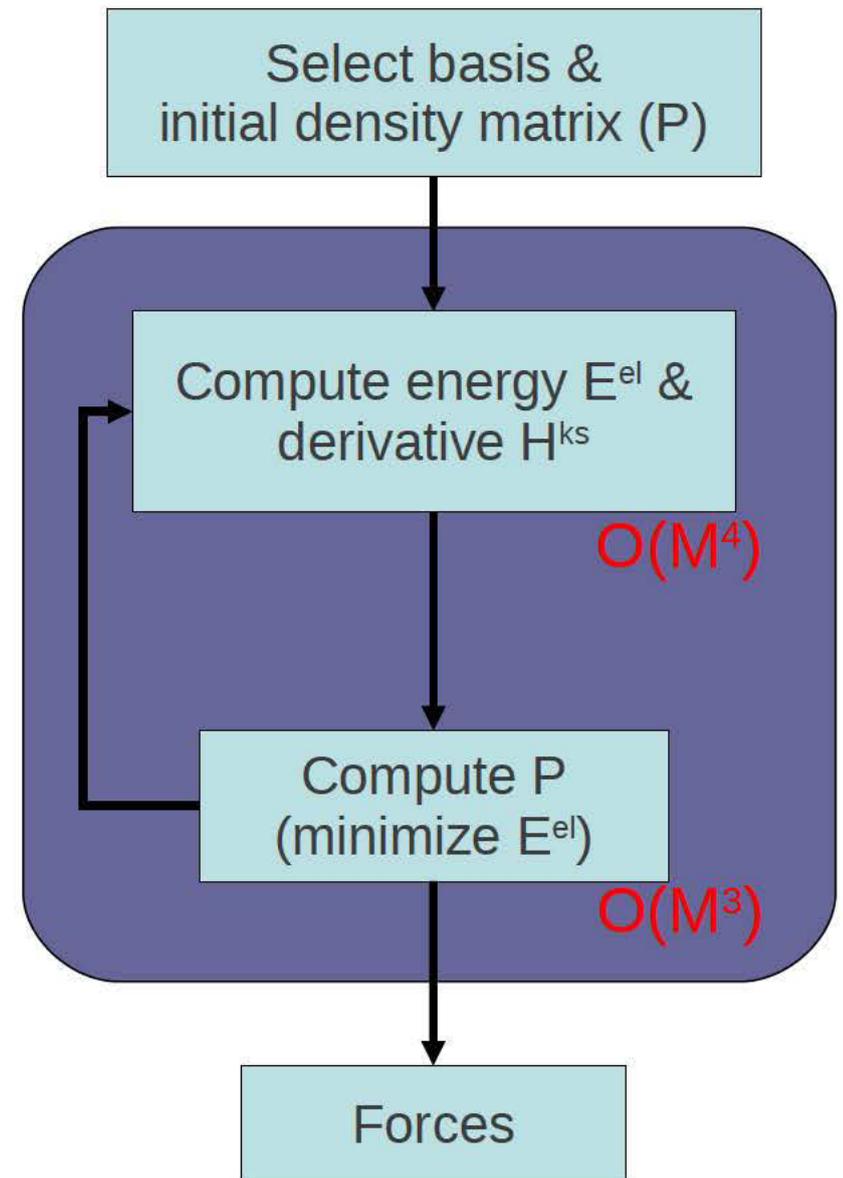
$$E^{el}[P^{\mu\nu}] = \sum_{\mu\nu} P^{\mu\nu} \int \varphi_{\mu}(r) \left(-\frac{\Delta}{2}\right) \varphi_{\nu}(r)$$

$$+ \sum_{\mu\nu} P^{\mu\nu} \int \int \varphi_{\mu}(r) V_{sep}^{PP}(r, r') \varphi_{\nu}(r')$$

$$+ \frac{1}{2} \int \int \frac{n(r)n(r')}{|r-r'|}$$

$$+ \int n(r) \varepsilon_{xc}[n](r)$$

Formally  $O(M^4) \rightarrow O(M)$  with GPW



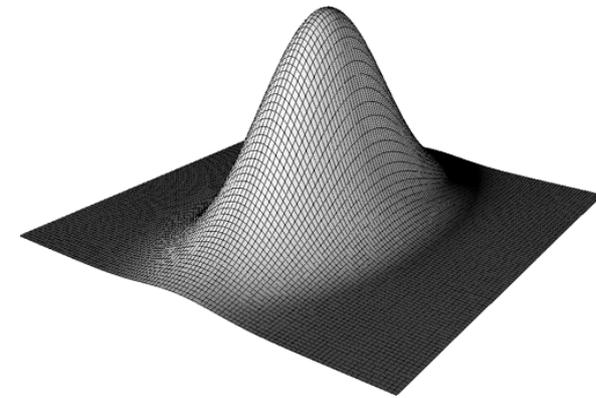
# A closer look at GPW

Real space (rs) density mapping and integration  
 Fourier space (FFT) for the coulomb problem

Only a linear number of terms have a non-zero contribution to the sum.  
 Each term affects only a constant volume in space

$O(N)$  ↗

$$\sum_{\mu\nu} P^{\mu\nu} \underbrace{\varphi_{\mu}(r)}_{RS} \underbrace{\varphi_{\nu}(r)}_{RS} \underbrace{\Rightarrow n(r)}_{RS} \underbrace{\Rightarrow n(G)}_{FFT}$$



$$\Rightarrow V_H(G) = \frac{4\pi n(G)}{G^2}, \quad E_H = \Omega \sum_G n^{cc}(G) V_H(G) \underbrace{\Rightarrow V_H(r)}_{FFT}$$

$$\underbrace{\Rightarrow V_{\mu\nu}}_{RS} = \int V_H(r) \varphi_{\mu}(r) \varphi_{\nu}(r)$$

A Gaussian basis allows for a very efficient procedure to compute the density on the grid

# Orbital transformations (OT)

A cubic, very robust algorithm avoiding the of traditional diagonalization

- New variables

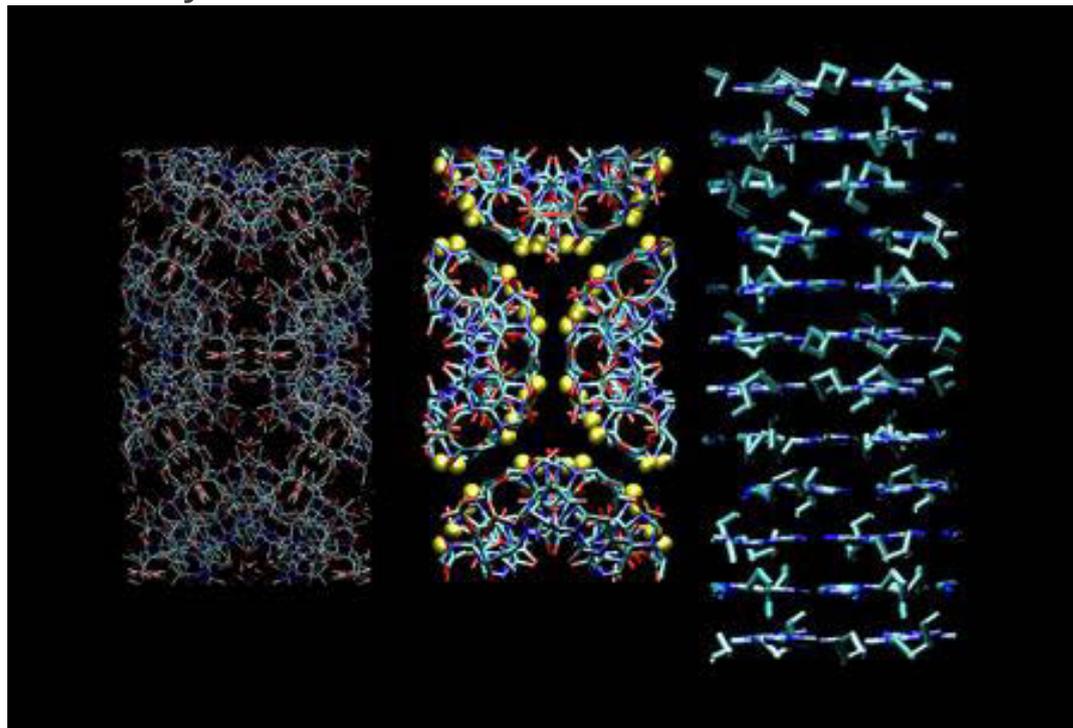
$$C(X) = C_0 \cos(\sqrt{X^T S X}) + X \frac{\sin(\sqrt{X^T S X})}{\sqrt{X^T S X}}$$

$$X^T S C_0 = 0 \quad C(X)^T S C(X) = 1 \quad \forall X$$

- Direct minimization of  $E_{\text{KS}}[\{X\}]$
- Linear constraint -> guaranteed convergence!

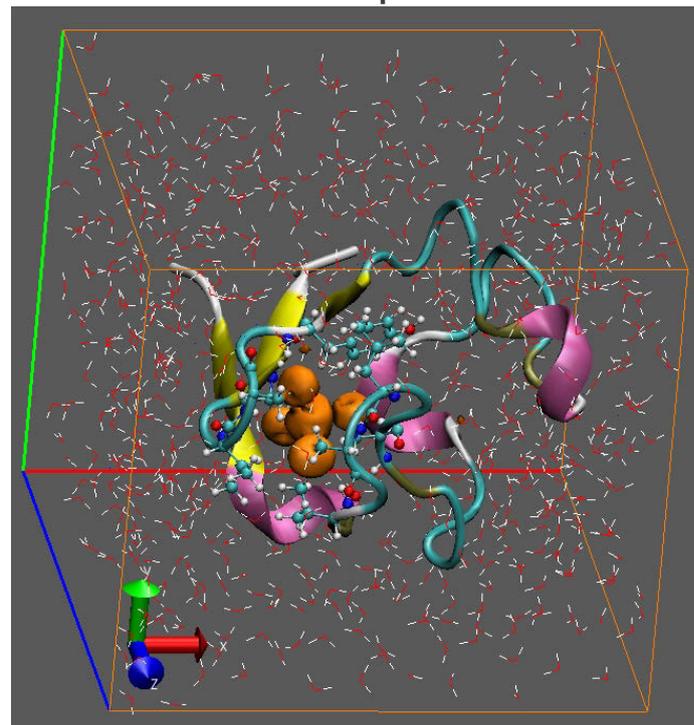
# DFT for large systems

DNA crystal



2388 atoms

Solvated metallo-protein



2825 atoms

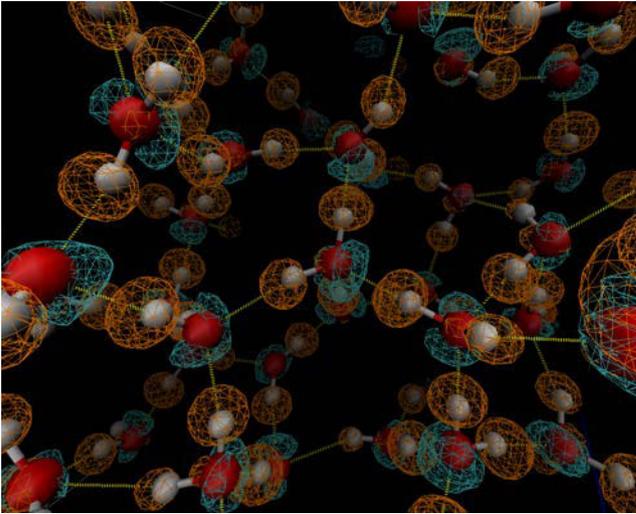
Linear scaling construction of the Kohn-Sham matrix,  
robust and efficient  $O(N^3)$  electronic minimization  
Parameter free & out-of-the-box for H-Rn

J. VandeVondele, J. Hutter, JCP 118, 4365-4369 (2003).

Sulpizi, M.; Raugei, S.; VandeVondele, J.; Carloni, P.; Sprik, M. JPCB 111, 3969 (2007).

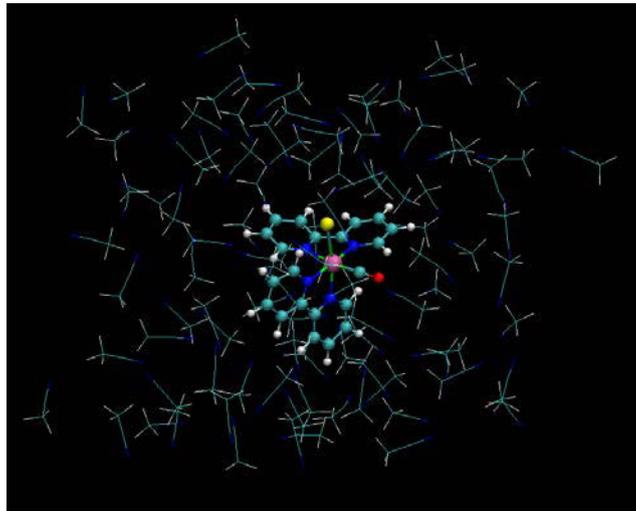
# Robust Ab initio MD

'Simple' liquids



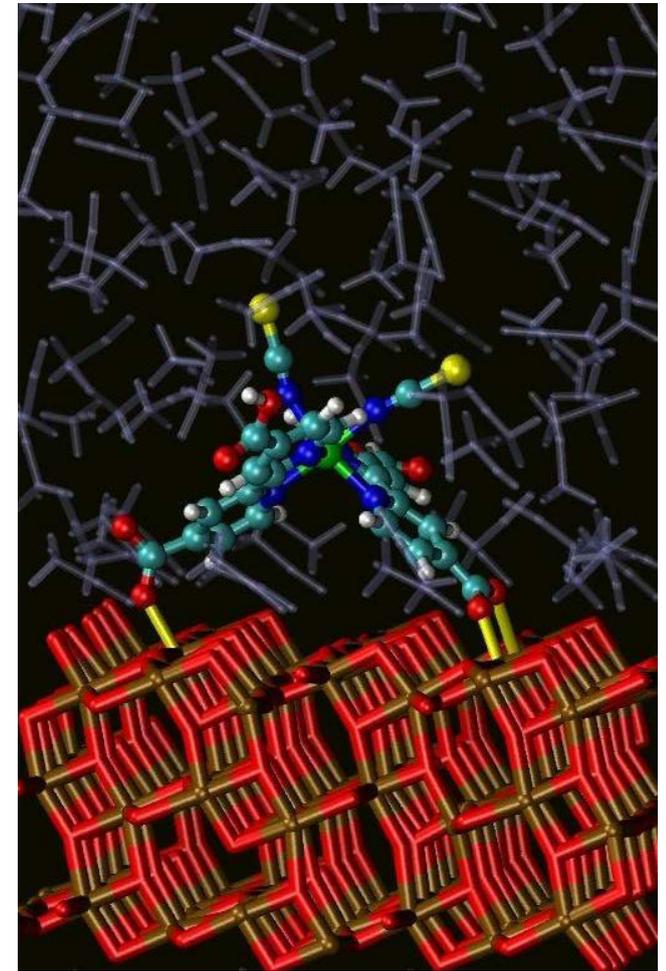
Hydrogen bonding in H<sub>2</sub>O

Solutes in explicit solvent



Ru(bpy)<sub>2</sub>COCl in acetonitrile,  
[21.43Å]<sup>3</sup> or 620 Atoms  
e.g. Redox properties

Complex Interfaces

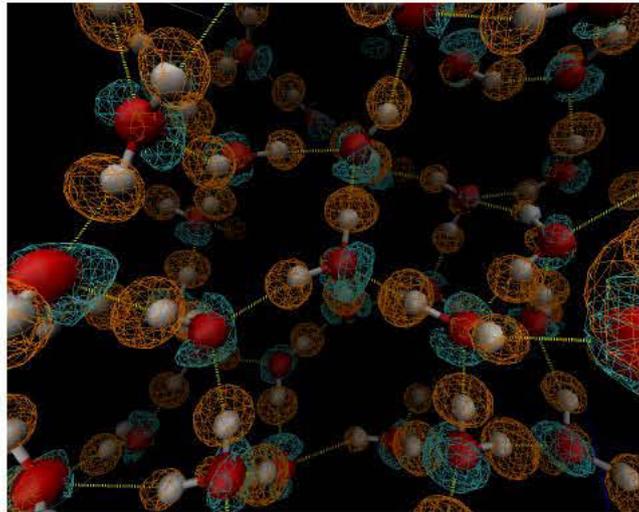


In situ IR spectroscopy  
(1300 atoms)

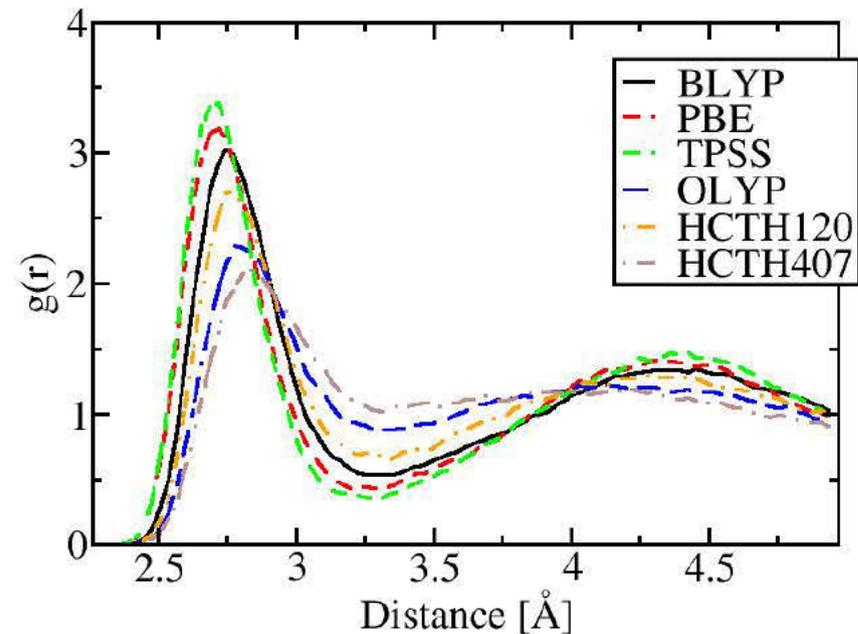
# Water: the most important liquid

In depth investigation of the performance of DFT for liquid water

Hybrid functionals

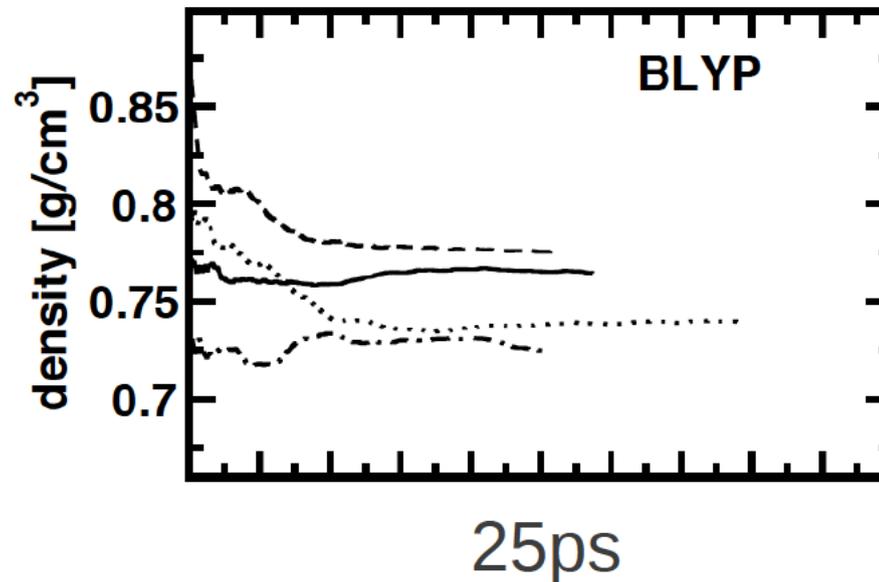


local functionals



# Water: the most important liquid

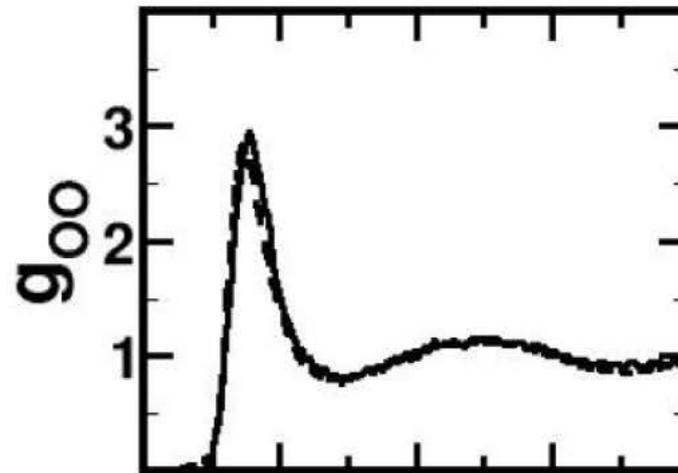
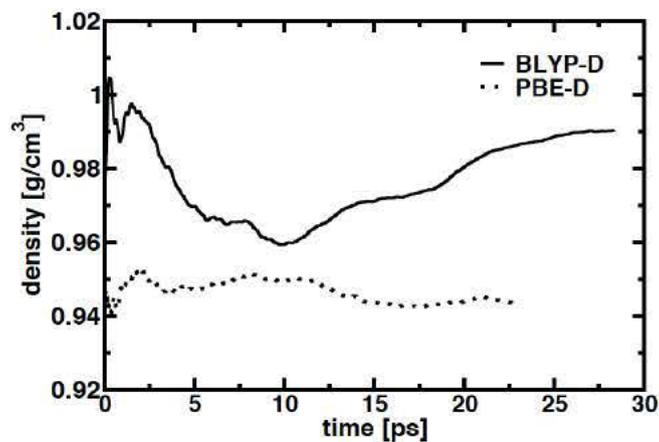
Simulation in the NPT ensemble



Standard functional gets density very wrong

# Water: the most important liquid

Simulation in the NPT ensemble



- Dispersion corrected BLYP (BLYP-D) yields:
- correct density
  - excellent pair correlation function

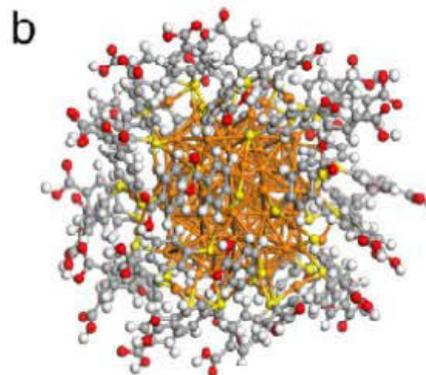
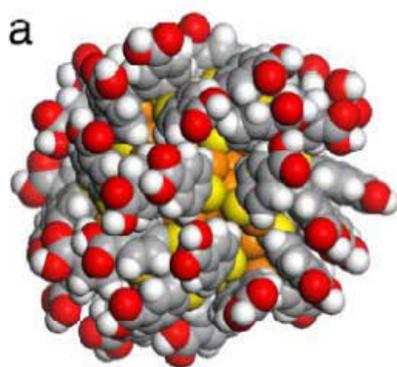
Generally: D correction improves energetics for biologically relevant interactions, such as pi-stacking and hydrogen bonding. Is empirical, but computationally free.

# CP2K: science (I)

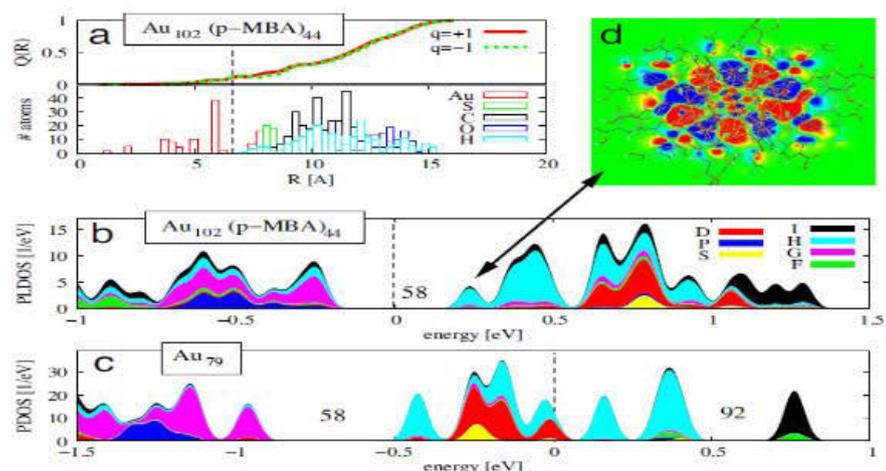
## A unified view of ligand-protected gold clusters as superatom complexes

Michael Walter<sup>†</sup>, Jaakko Akola<sup>††</sup>, Olga Lopez-Acevedo<sup>†</sup>, Pablo D. Jadzinsky<sup>5††</sup>, Guillermo Calero<sup>5</sup>, Christopher J. Ackerson<sup>5||</sup>, Robert L. Whetten<sup>††</sup>, Henrik Grönbeck<sup>††</sup>, and Hannu Häkkinen<sup>†55††||</sup>

PNAS July 8, 2008 vol. 105 no. 27 9157–9162



System size: 762 Atoms  
~3400 electrons



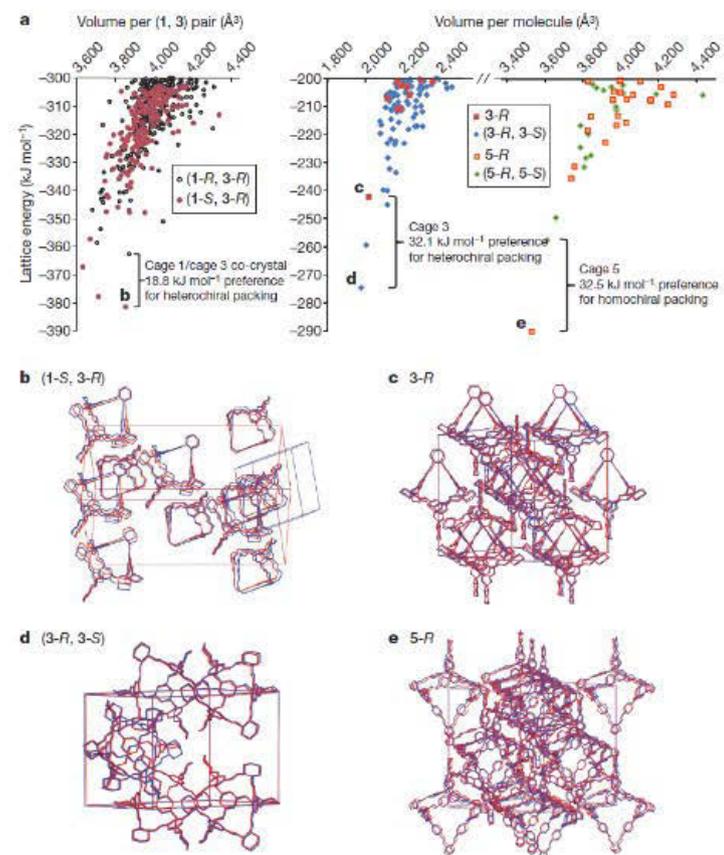
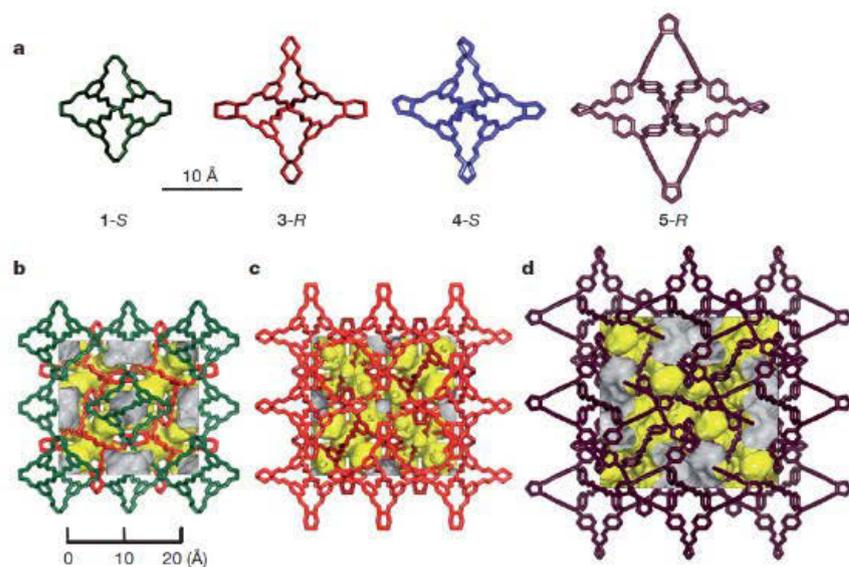
Electronic structure of nanoparticles

# CP2K: science (II)

## Modular and predictable assembly of porous organic molecular crystals

James T. A. Jones<sup>1</sup>, Tom Hasell<sup>1</sup>, Xiaofeng Wu<sup>1</sup>, John Bacsá<sup>1</sup>, Kim E. Jelfs<sup>1</sup>, Marc Schmidtman<sup>1</sup>, Samantha Y. Chong<sup>1</sup>, Dave J. Adams<sup>1</sup>, Abbie Trewin<sup>1</sup>, Florian Schiffman<sup>2</sup>, Furio Cora<sup>2</sup>, Ben Slater<sup>2</sup>, Alexander Steiner<sup>1</sup>, Graeme M. Day<sup>3</sup> & Andrew I. Cooper<sup>1</sup>

16 JUNE 2011 | VOL 474 | NATURE | 367



Structure prediction of metal organic frameworks

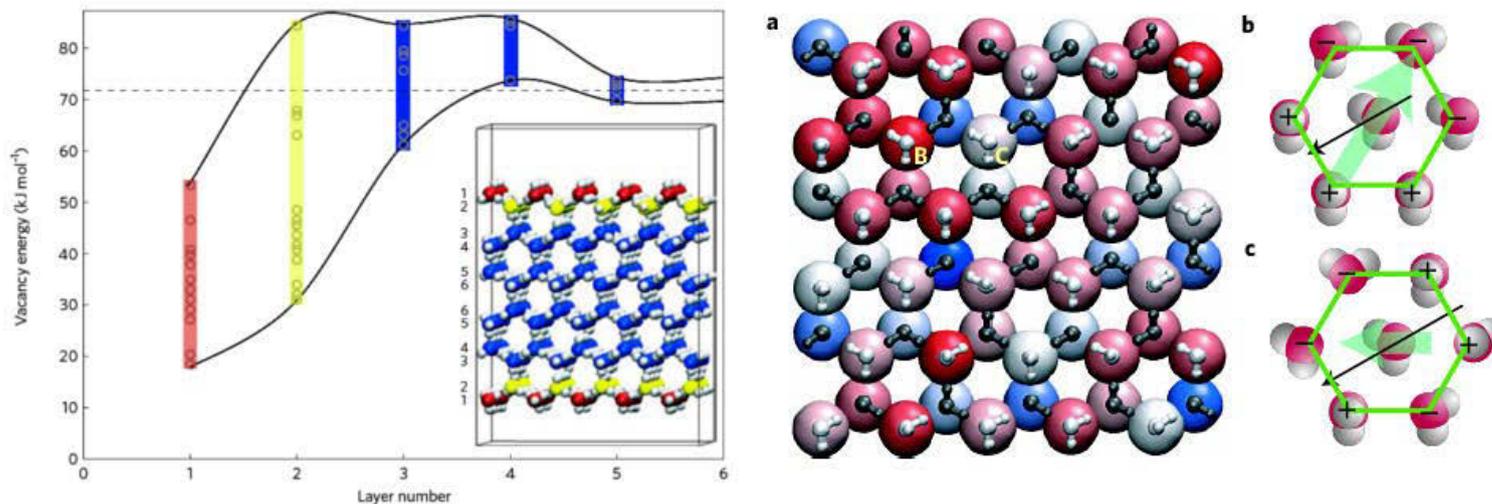
# CP2K: science (III)

## Large variation of vacancy formation energies in the surface of crystalline ice

M. Watkins<sup>1,2,3</sup>, D. Pan<sup>4</sup>, E. G. Wang<sup>5</sup>, A. Michaelides<sup>1,2,3</sup>, J. VandeVondele<sup>6</sup> and B. Slater<sup>1,3\*</sup>

<sup>1</sup>Department of Chemistry, Christopher Ingold Building, 20 Gordon Street, University College London, London WC1H 0AJ, UK, <sup>2</sup>London Centre for Nanotechnology, University College London, London WC1H 0AJ, UK, <sup>3</sup>TYC@UCL, University College London, London WC1H 0AJ, UK, <sup>4</sup>Institute of Physics, Chinese Academy of Sciences, PO Box 603, Beijing 100190, China, <sup>5</sup>School of Physics, Peking University, Beijing 100871, China, <sup>6</sup>Institute of Physical Chemistry, University of Zurich, Winterthurerstrasse 190, CH-8057 Zurich, Switzerland. \*e-mail: b.slater@ucl.ac.uk.

NATURE MATERIALS | VOL 10 | OCTOBER 2011



Disordered and frustrated materials

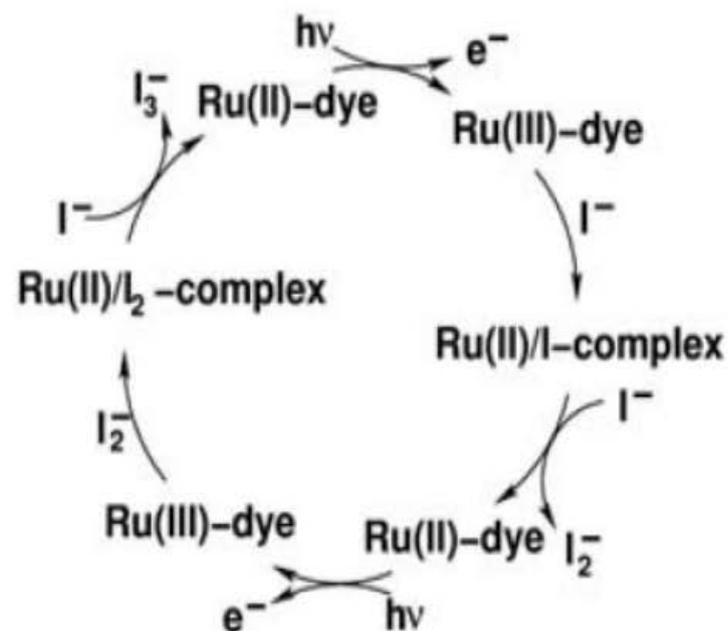
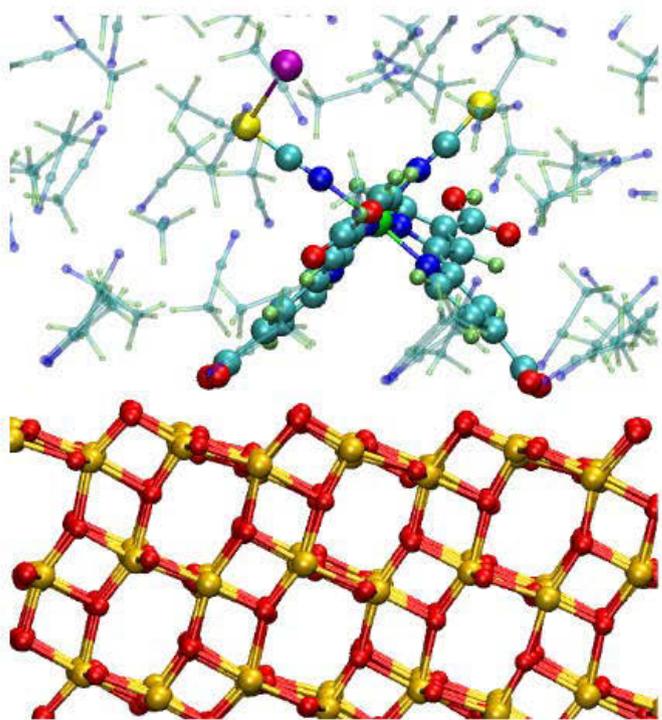
# CP2K: science (IV)

## An atomistic picture of the regeneration process in dye sensitized solar cells

Florian Schiffmann<sup>a</sup>, Joost VandeVondele<sup>a,1</sup>, Jürg Hutter<sup>a</sup>, Atsushi Urakawa<sup>b</sup>, Ronny Wirz<sup>b</sup>, and Alfons Baiker<sup>b</sup>

<sup>a</sup>Institute of Physical Chemistry, University of Zurich, Winterthurerstrasse 190, 8057 Zurich, Switzerland; and <sup>b</sup>Department of Chemistry and Applied Biosciences, Institute for Chemical and Bioengineering, ETH Zurich, Hönggerberg, HCI, 8093 Zurich, Switzerland

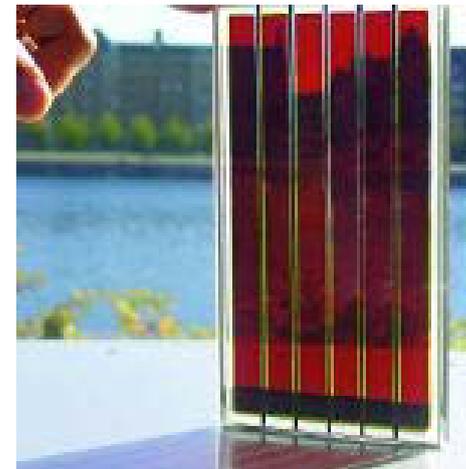
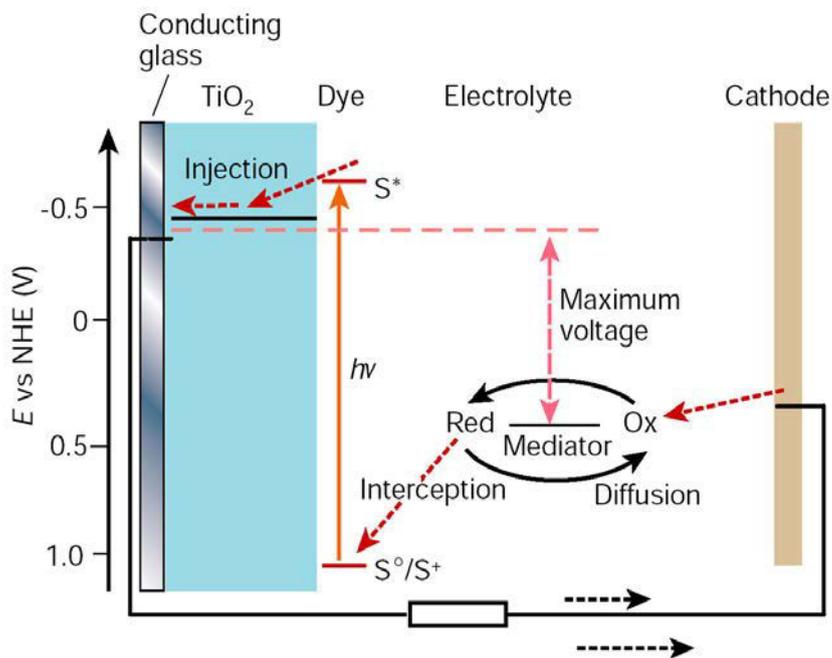
4830–4833 | PNAS | March 16, 2010 | vol. 107 | no. 11



Functionalized solid/liquid interfaces

# Dye sensitized solar cells (DSSC)

Sustainable energy production converting sunlight into electricity



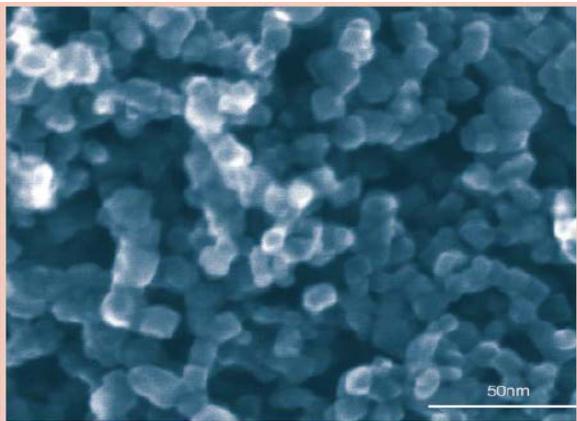
Grätzel, Nature (1991,2001)

~10% efficiency, short energy-payback, tolerant to impurities, inexpensive technologies, cheap materials, wide temperature range, diffuse light OK, various colors (semi-transparent), flexible

# DSSC: atomistic engineering

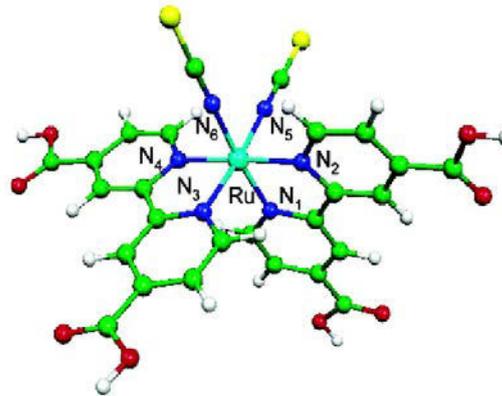
Prototypical high performance cells:

Mesoporous  $\text{TiO}_2$  (Anatase)



Grätzel, Nature (1991,2001)

Ruthenium Dye (N3)

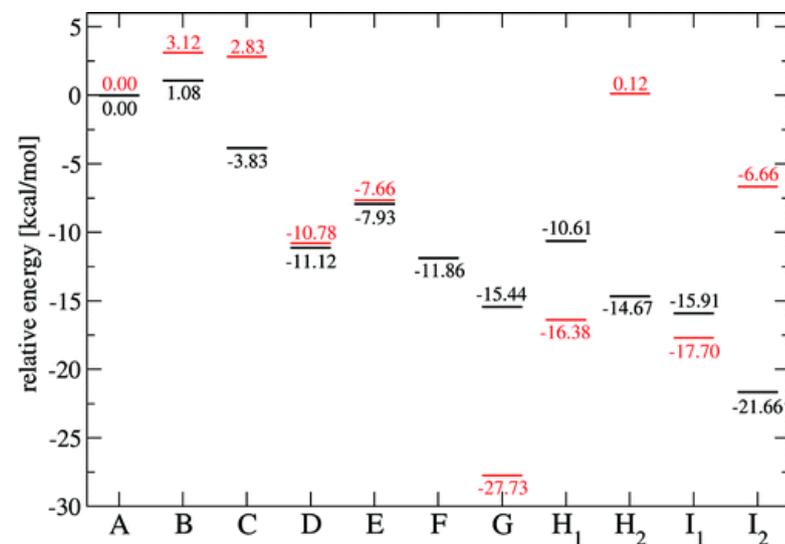
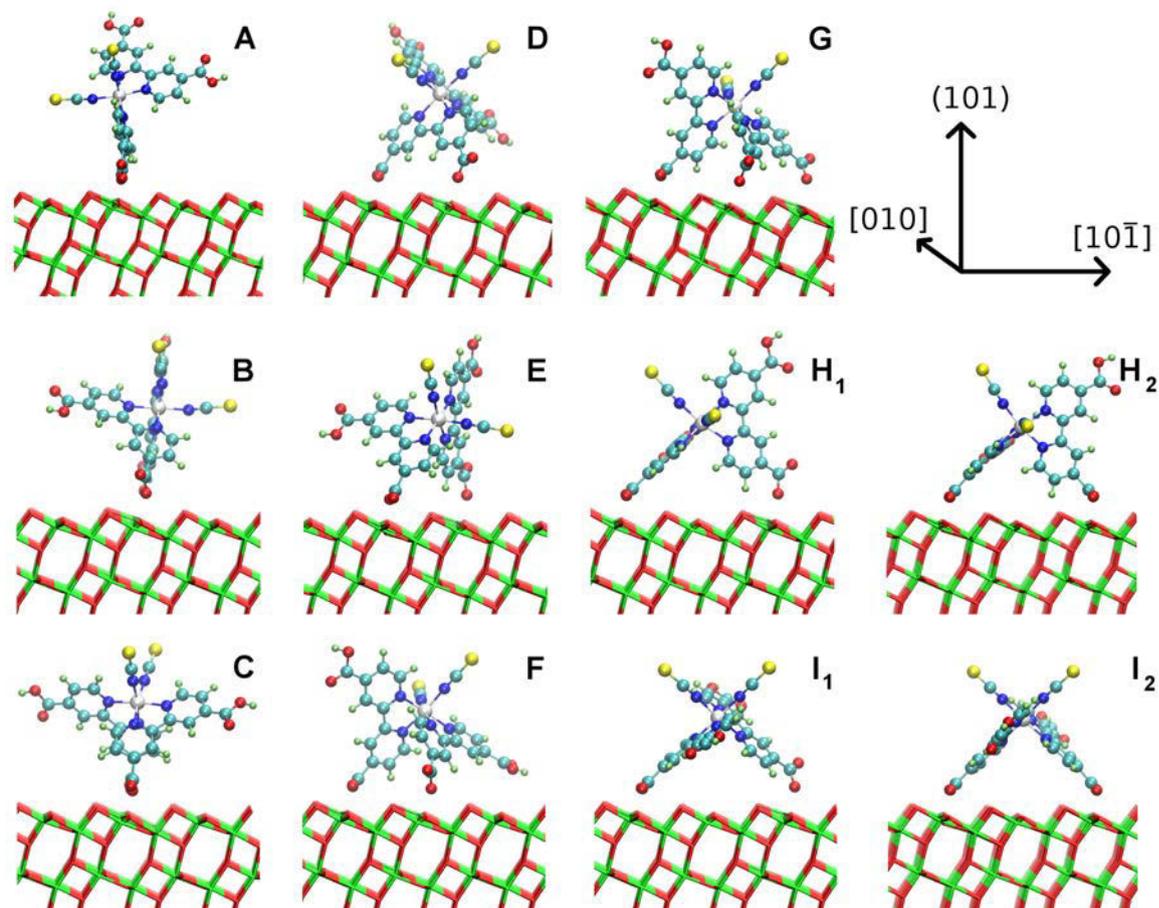


$\text{I}^-/\text{I}_3^-$  in Acetonitrile



Taylor the interface to improve efficiency, stability and cost

# N3 binding on anatase(101)

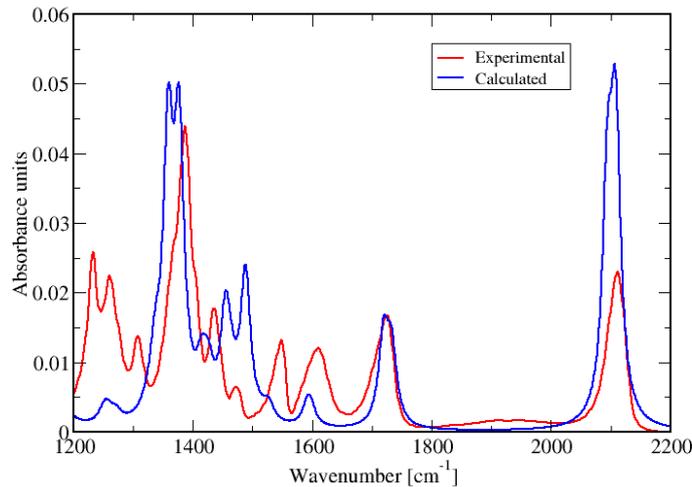


Relative binding energies are indicative:

- 2-3 carboxylate groups interact
- 2 bpy ligands involved
- Monodentate and bridged binding possible
- Protonation strongly influences relative stability

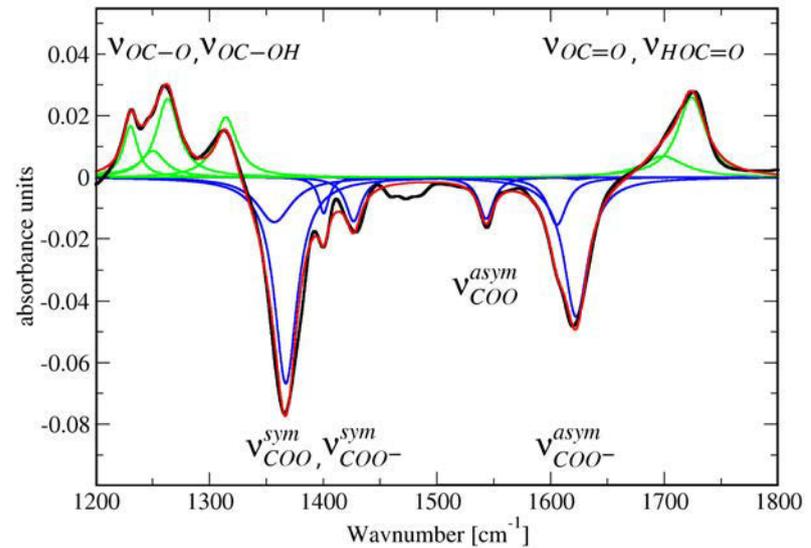
# IR spectra: Theory & Experiment

Theory & Experiment



Promising agreement ( $\pm 30 \text{ cm}^{-1}$ ),  
a unique binding geometry is not identified.

Exp. difference spectrum N3/N712

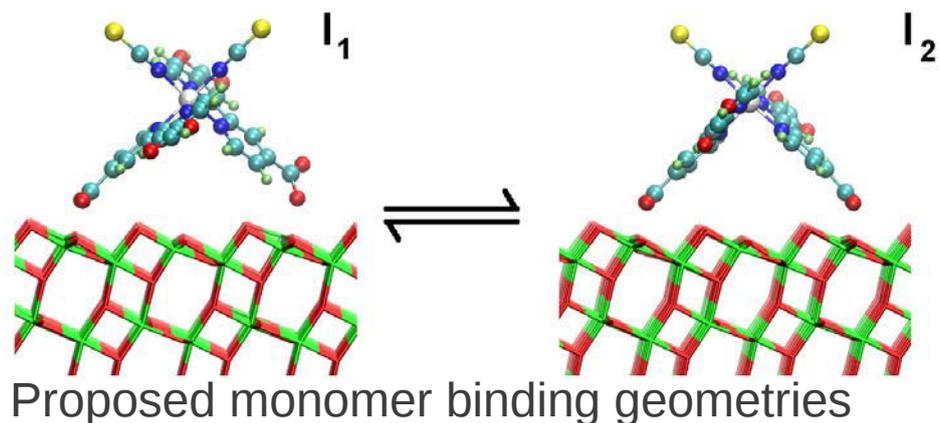


Theory provides assignments

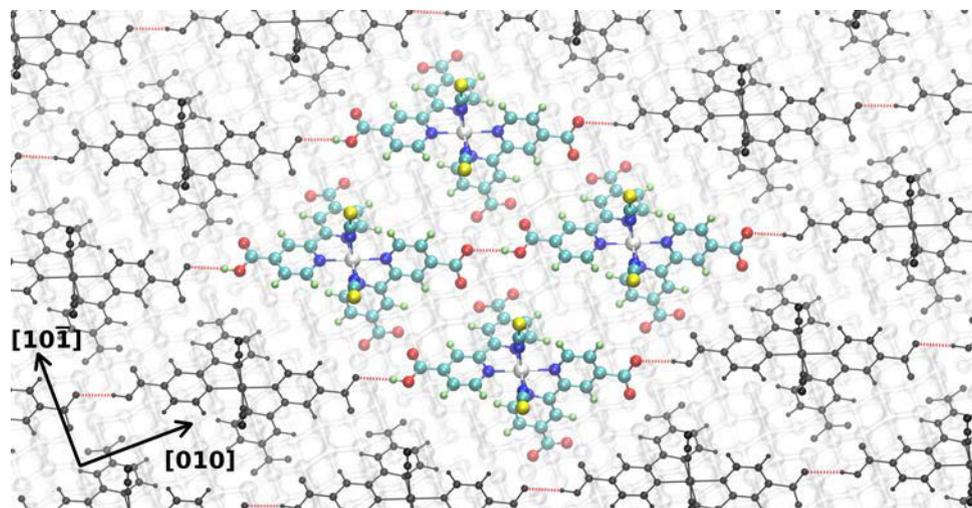
## Conclusions:

- base-induced desorption of dye adsorbed in the presence of protons
- change in binding mode depending on protonation

# Suggested binding mode



Equilibrium between two Configurations (pH dependent)

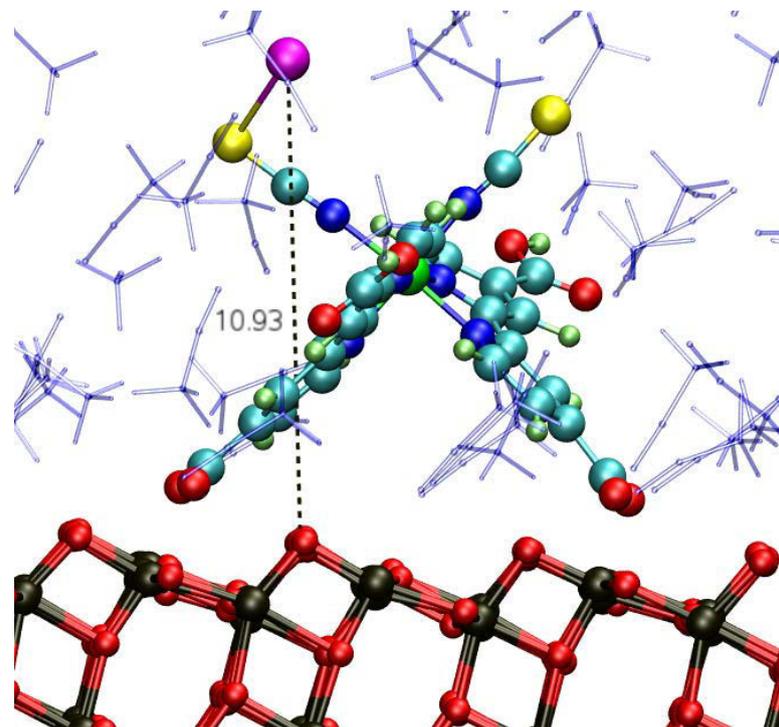


Computed self-assembly of  $I_2$  on anatase(101)

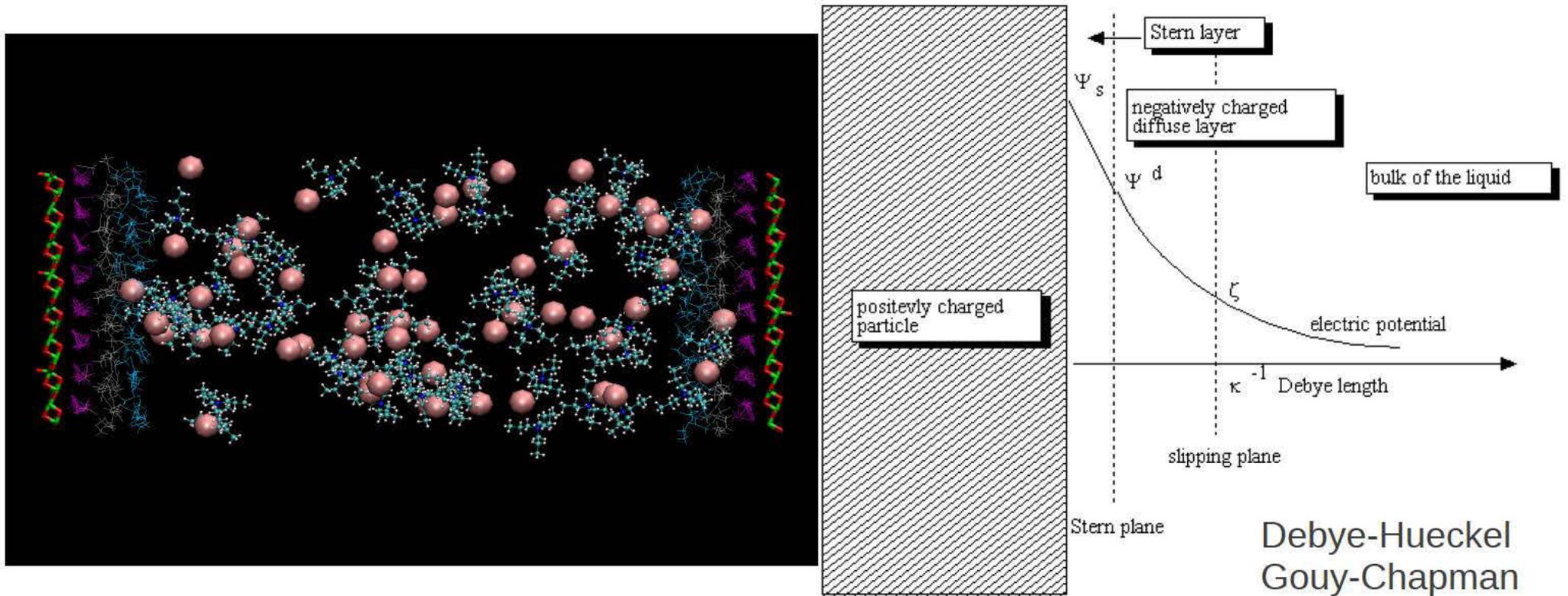
Experimental data consistent with self-assembly:

- Dense packing
- Dimers in STM (rutile)
- Desorption with base

# Including the electrolyte



# Electrolyte near the interface



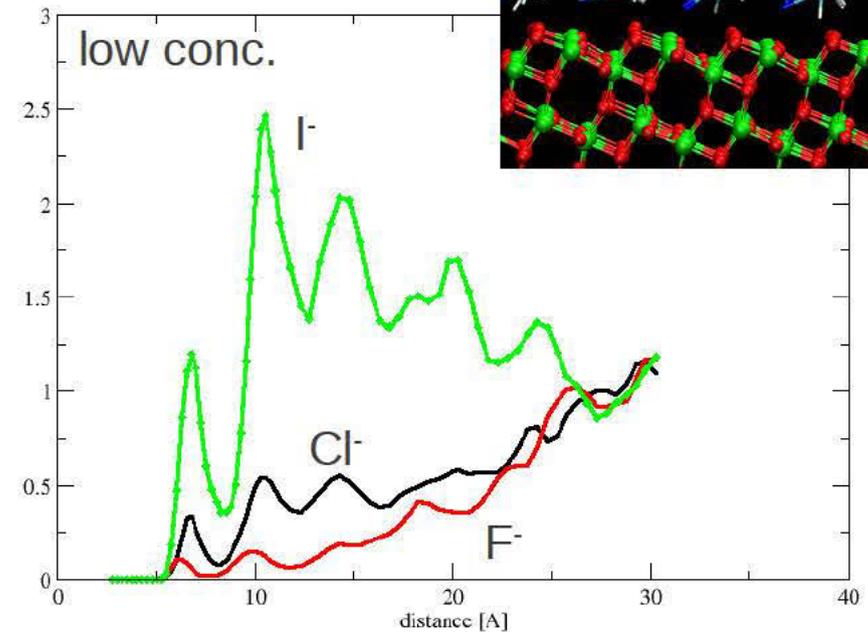
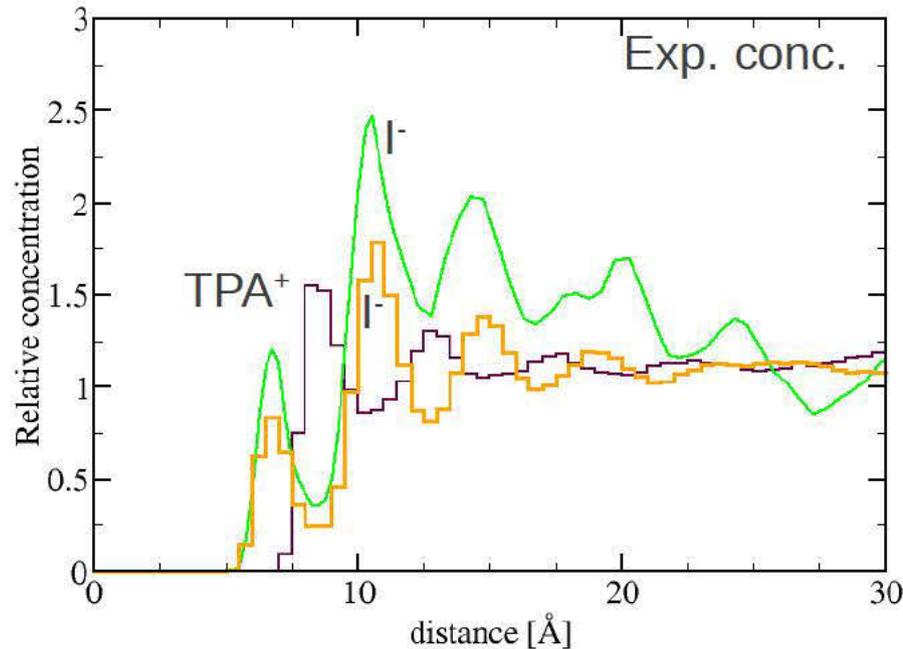
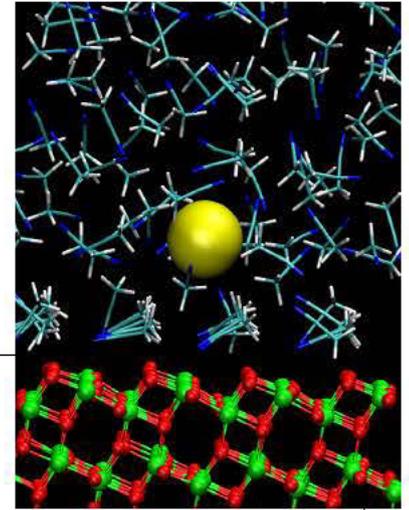
Atomistic

vs.

Continuum

# I<sup>-</sup> concentration profiles

- I<sup>-</sup> specific enhancement near the interface
- The first solvent layer is passivating the interface



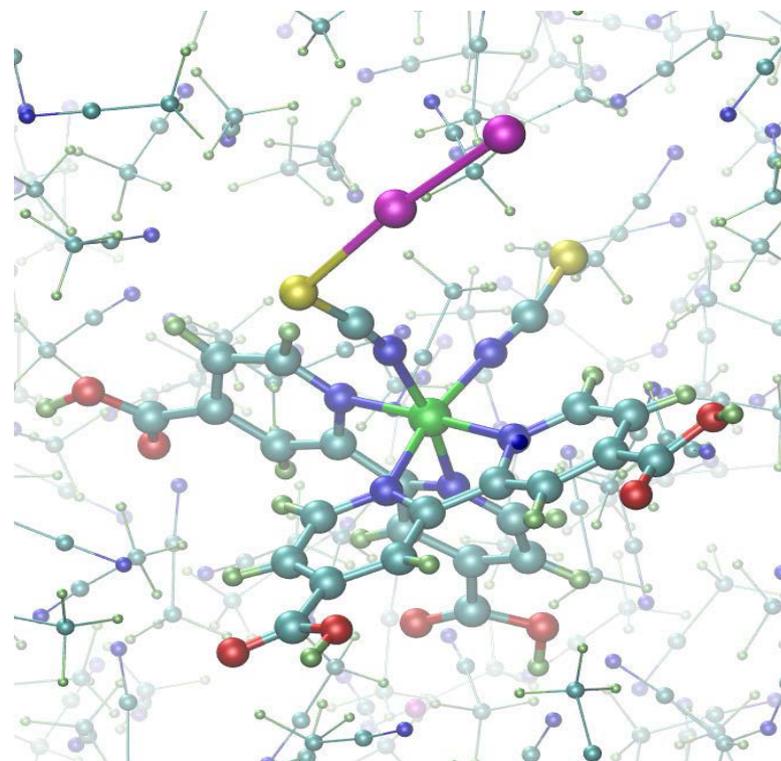
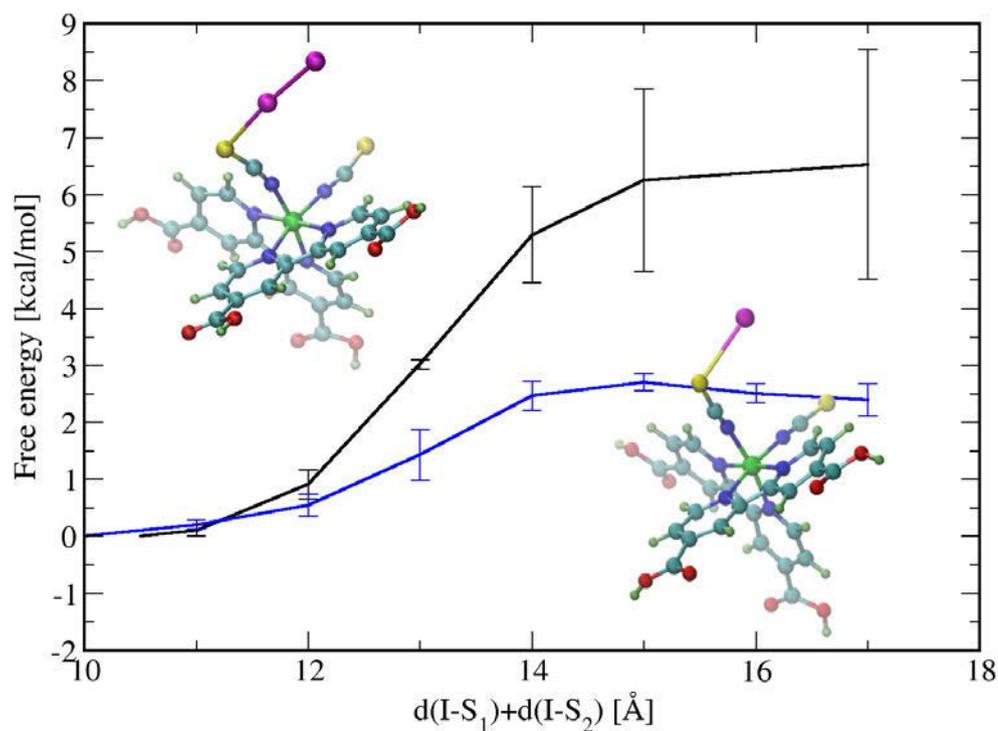
~microsecond simulation time (classical MD)

The molecular structure of the liquid near the interface is required to obtain a qualitatively correct ion distribution

# Dye - Iodide interaction

DFT calculations of complex formation.

Free energies of binding in explicit solution.

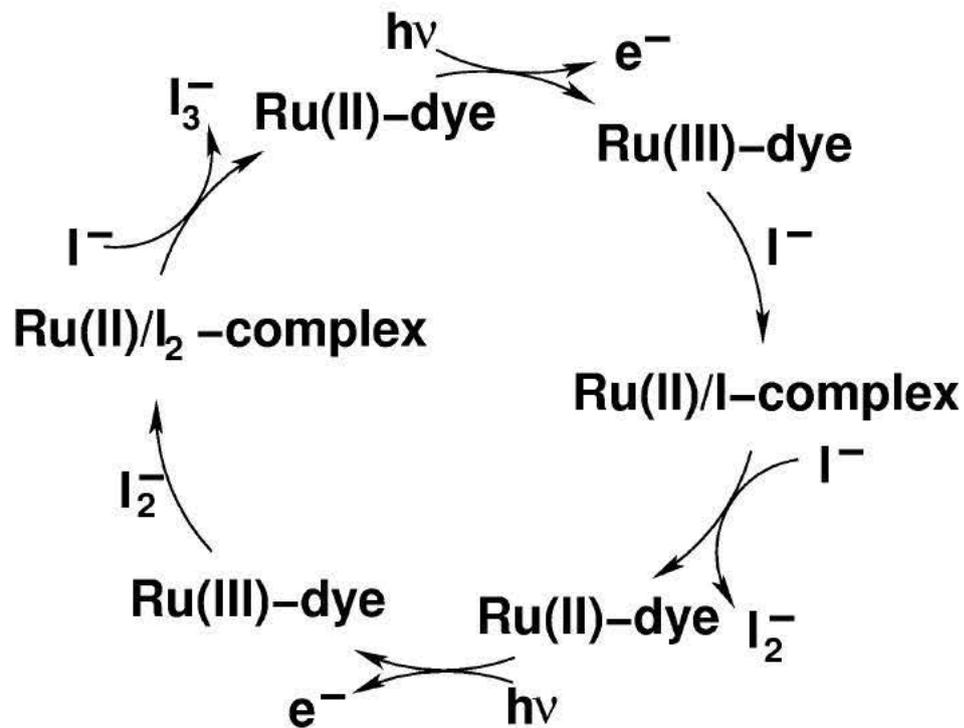
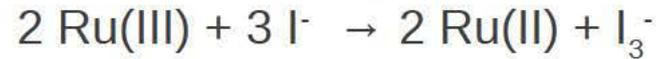


I<sup>-</sup>, I<sub>2</sub><sup>-</sup> form stable complexes in solution

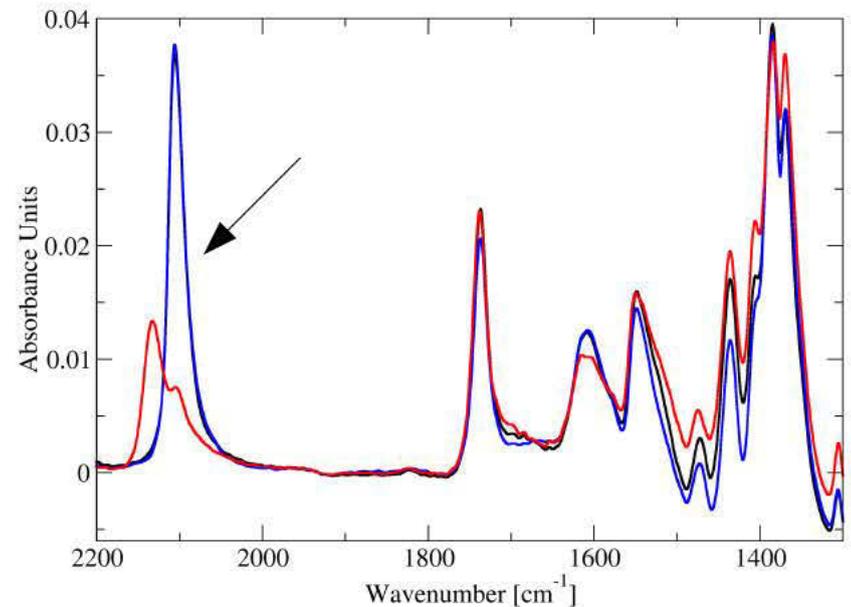
SCN<sup>-</sup> group essential in complex formation

# Dye regeneration

non-trivial:

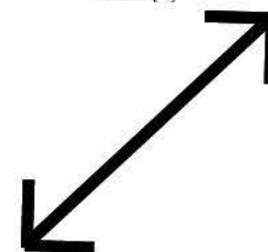
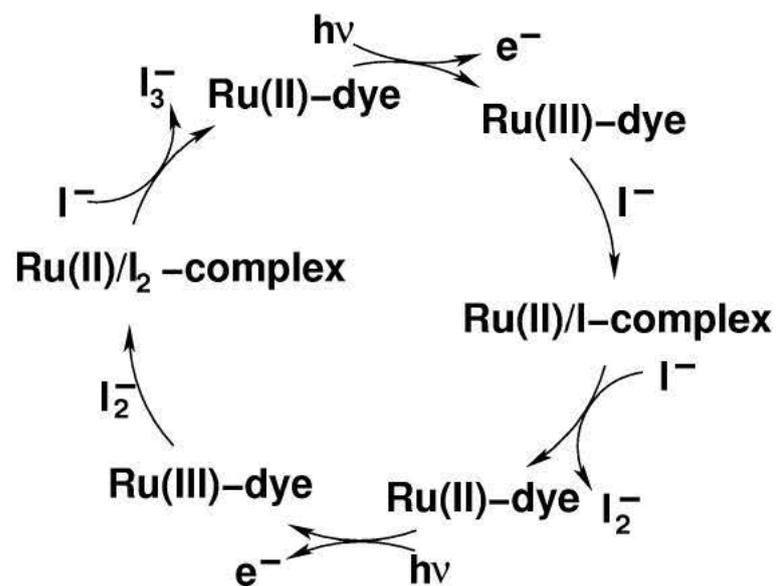
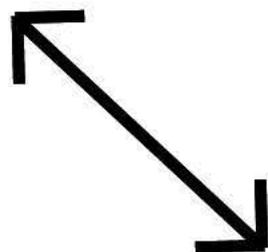
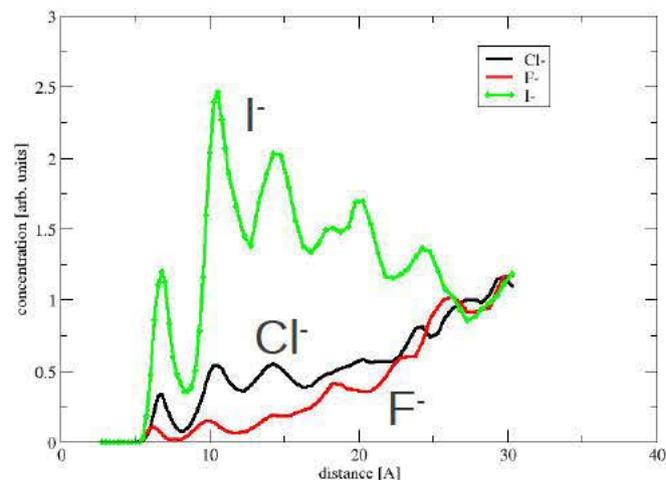
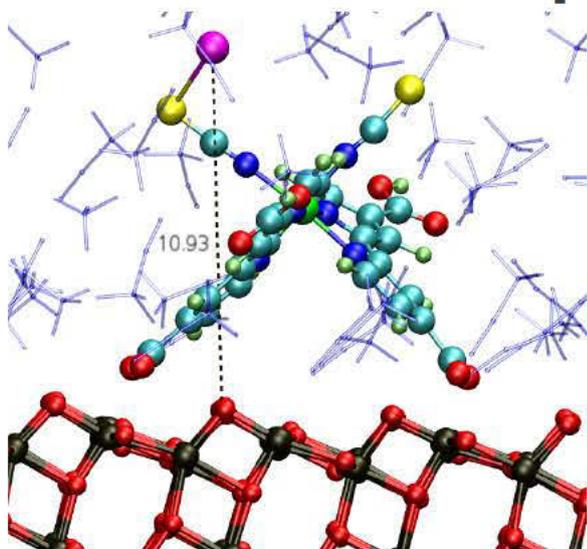


Exp. Confirmation of  $[\text{Dye:I}_2]$  complex



A highly efficient (barrierless) cycle.  
 $[\text{Dye:I}_2]$  and  $[\text{Dye:I}]$  complexes formed.

# Explanation of the superior performance



# Recent developments

Full linear scaling GGA DFT

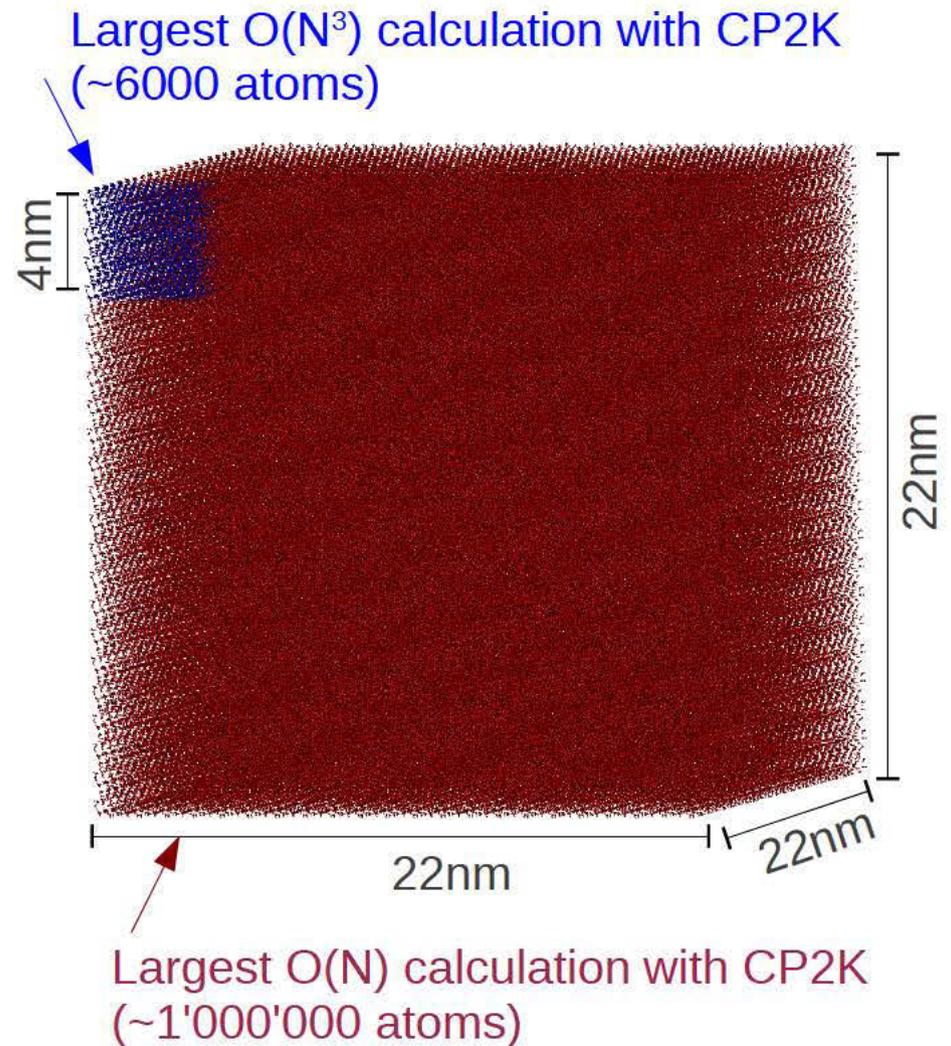
Efficiently using hybrid functionals

Linear scaling GGA DFT

# Linear Scaling SCF

Traditional approaches to solve the self-consistent field (SCF) equations are  $O(N^3)$  limiting system size significantly.

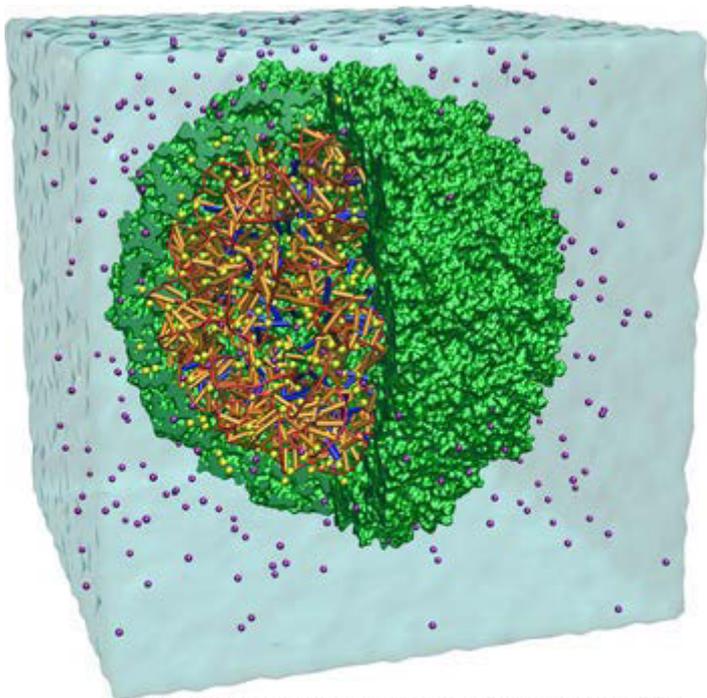
A newly implemented algorithm is  $O(N)$ , allowing for far larger systems to be studied.



# Linear Scaling SCF

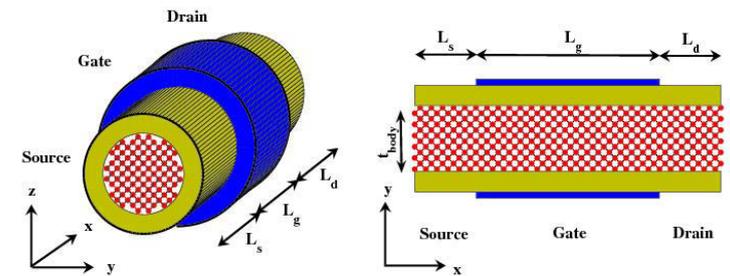
New regime: small devices, heterostructures, interfaces, nano-particles, a small virus.

Solvated STMV: 1M

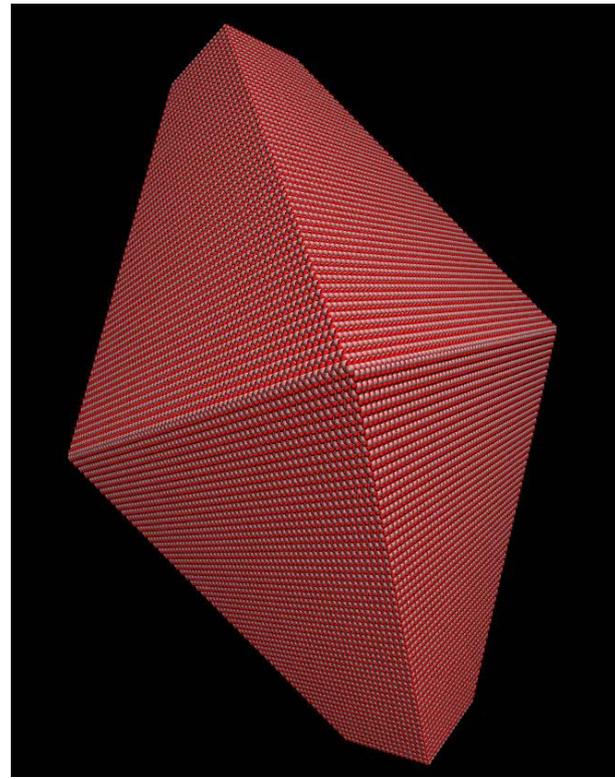


Theoretical and Computational Biophysics Group  
Beckman Institute  
University of Illinois at Urbana-Champaign

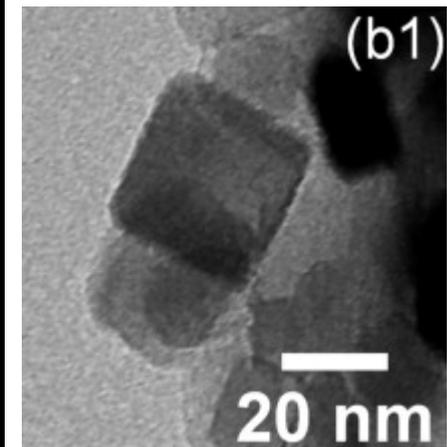
Gate-all-around FET



With Mathieu Luisier



1.5M atoms  
Anatase nanocrystal



Caplovicova et al.  
App. Cat. B, 224, 117

# Sign matrix iterations

The density matrix ( $P$ ) is function of  $H$

$$P = \frac{1}{2}(I - \text{sign}(S^{-1}H - \mu I))S^{-1}.$$

A simple iterative scheme (Newton-Schultz) gives  $\text{sign}(A)$ :

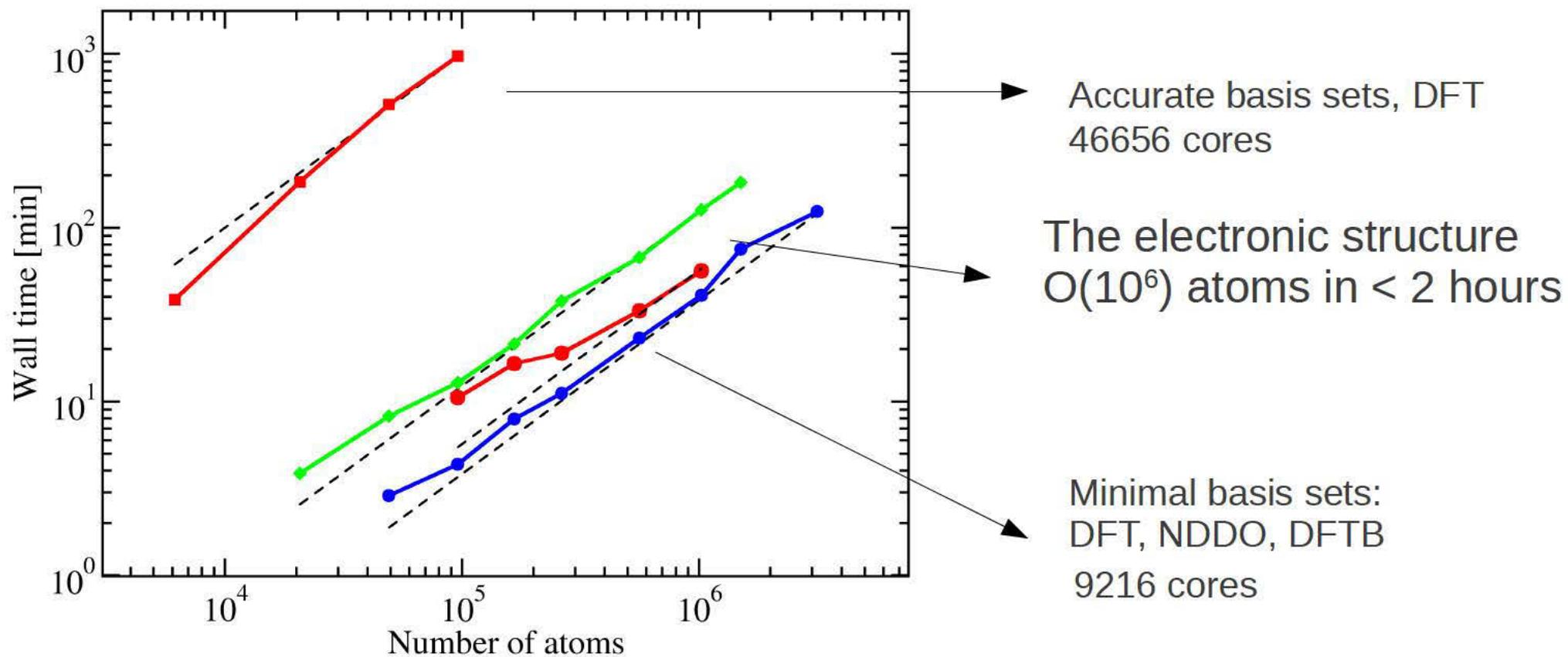
$$X_{n+1} = \frac{1}{2}X_n(3I - X_n^2).$$

Using only sparse matrix matrix multiplies (not SPMV!)  
linear scaling can be obtained



A dedicated sparse matrix multiply library is extremely important  
This library is being ported to GPUs

# Millions of atoms in the condensed phase

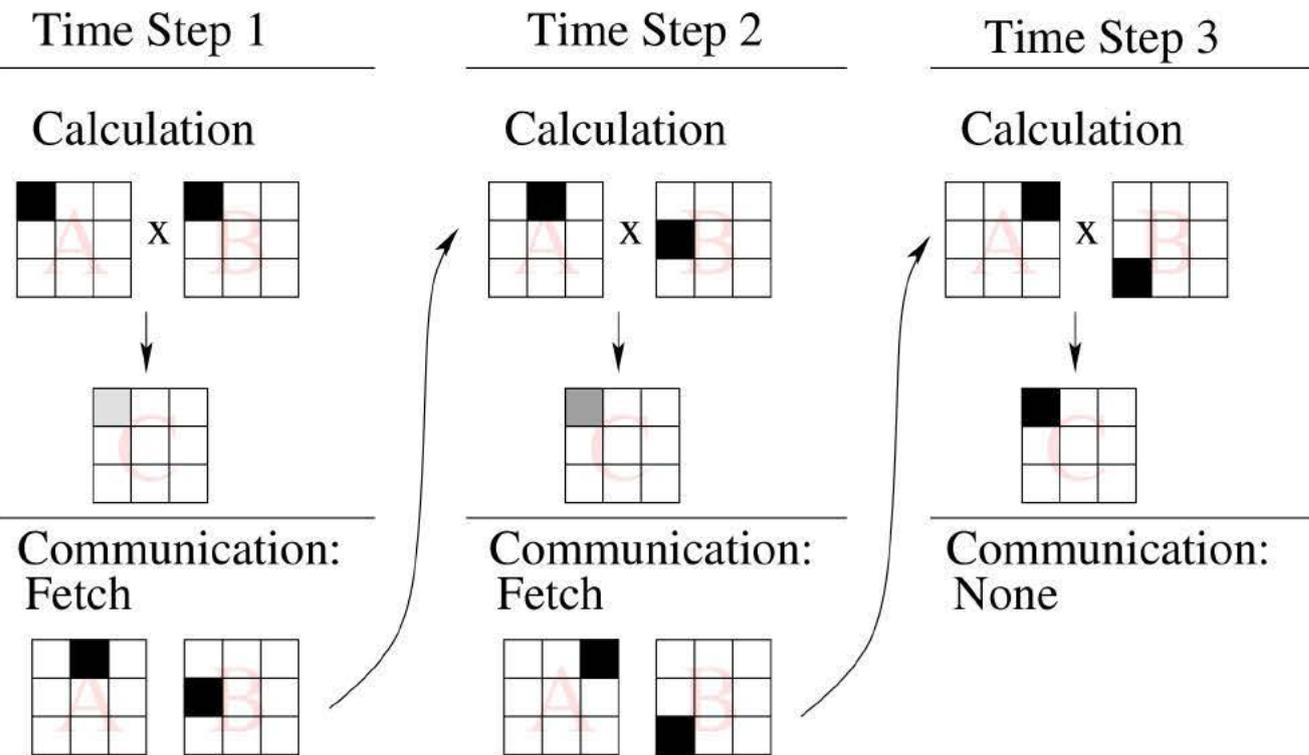


Bulk liquid water. Dashed lines represent ideal linear scaling.

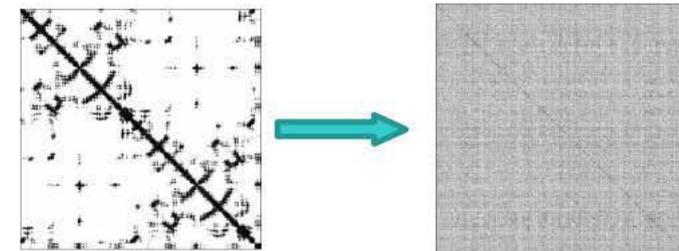
# DBCSR: a sparse matrix library

Distributed Blocked Compressed Sparse Row  
Distributed Blocked Cannon Sparse Recursive

Target the application: atoms → Blocks (e.g. 5x5, 13x13, 23x23)  
Linear scaling → Sparse  
Fully dense → Cannon  
Large scale → Distributed  
High Performance → Recursive

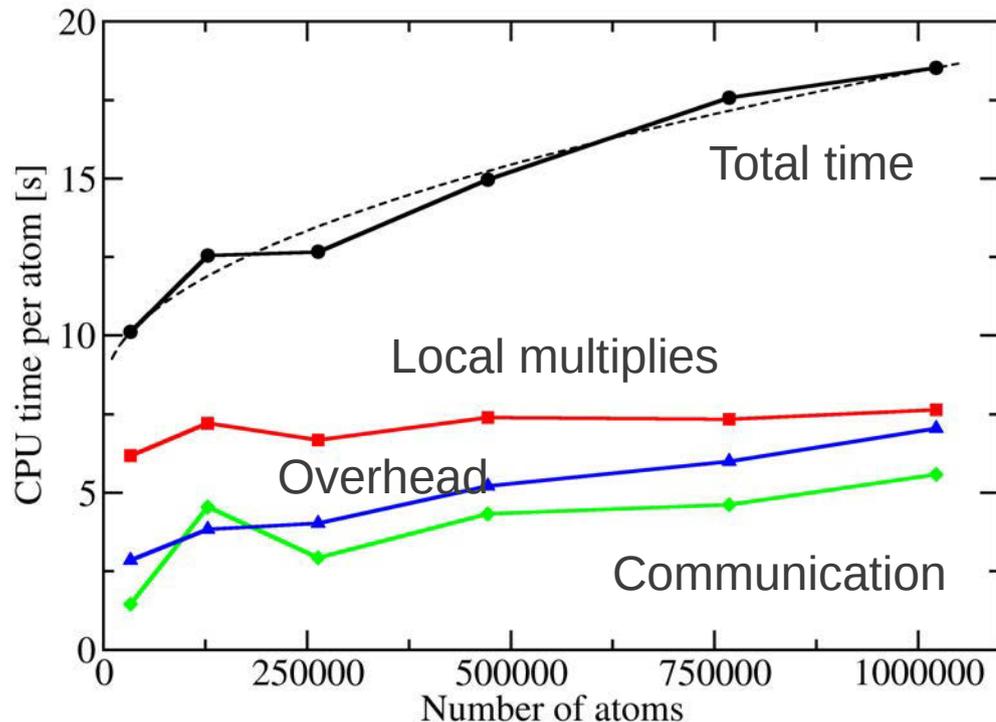


Cannon style communication on a homogenized matrix for strong scaling



# Towards $O(1)$ :

constant walltime with proportional resources



Stringent test:

Small blocks, large overhead

Very sparse matrices

Running with 200 atoms / MPI task

Local multiplies constant (OK!).

Overhead & Communication

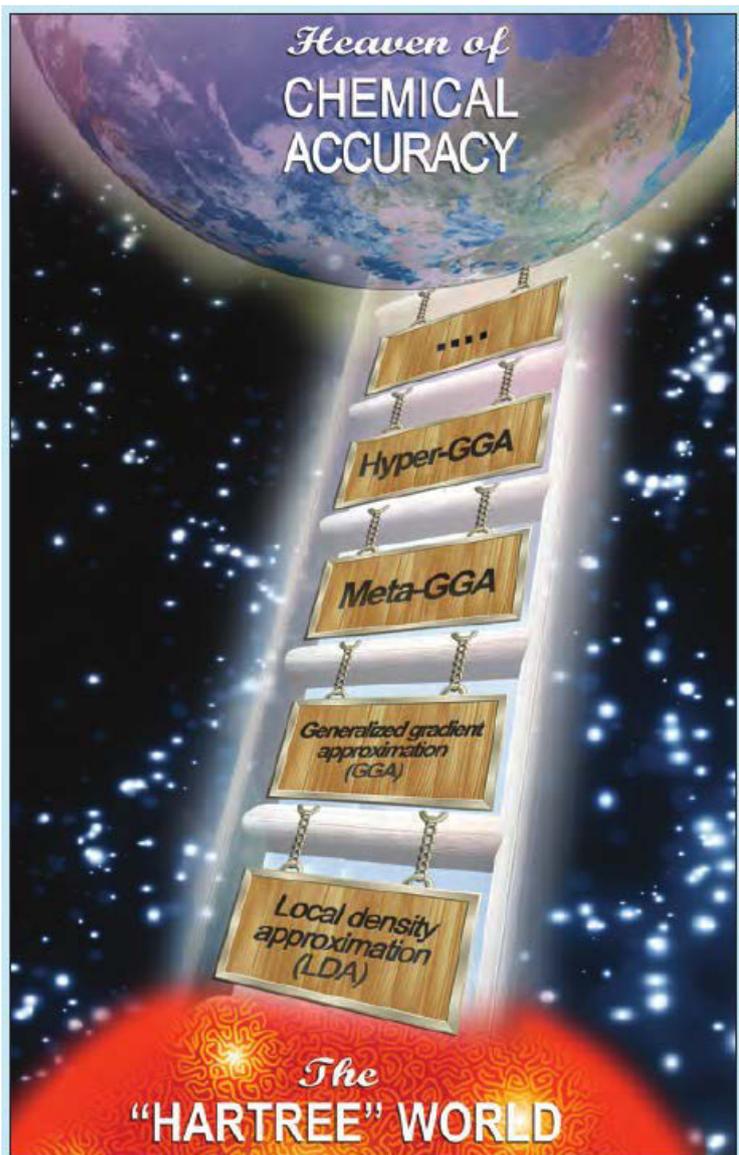
Grows with  $\sqrt{N}$

Needs a replacement for Cannon

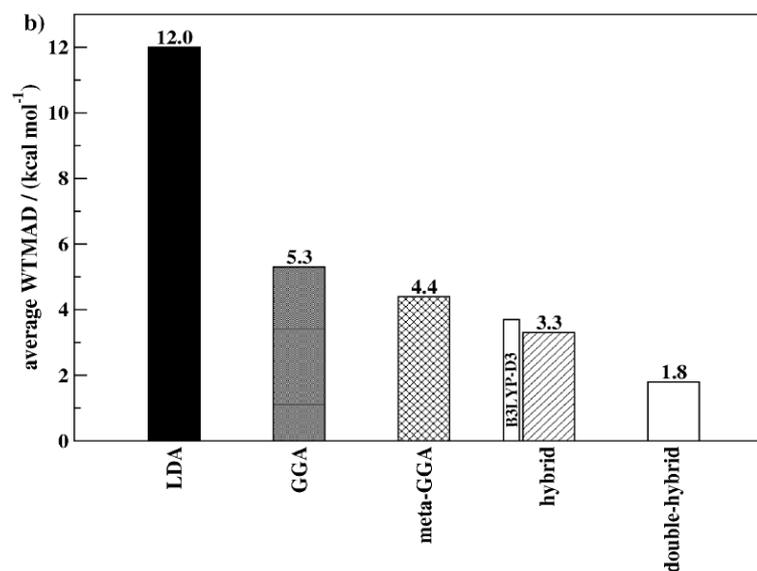
Work is underway to replace the Cannon algorithm with something new!  
Retain the  $\sqrt{N}$  max comm, yield constant comm in the limit.

# Hybrid functionals

# Advances in DFT



Exchange and correlation functionals of improving can be constructed by adding new ingredients:



Dispersion / van der Waals corrected functionals  
(Screened) Hybrid functionals

Goerigk, L. and Grimme, S.  
Phys. Chem. Chem. Phys. (2011) 13, 6670-6688.  
Mundy, Kathmann, Rousseau, Schenter,  
VandeVondele, Hutter, SCIDAC reviews (spring 2010).

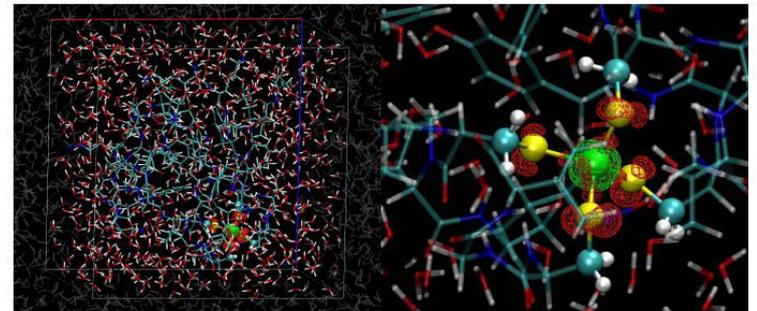
# Hartree-Fock exchange

$$E_x^{\text{HF}} = -\frac{1}{2} \sum_{\alpha\beta\gamma\delta} P_{\alpha\beta} P_{\gamma\delta} (\phi_{\alpha}\phi_{\gamma}|\phi_{\beta}\phi_{\delta})$$

$$(\phi_{\alpha}\phi_{\gamma}|\phi_{\beta}\phi_{\delta}) = \int d\mathbf{r}d\mathbf{r}' \frac{\phi_{\alpha}(\mathbf{r})\phi_{\gamma}(\mathbf{r})\phi_{\beta}(\mathbf{r}')\phi_{\delta}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

An easy term in Gaussian basis sets, but brute force scaling as  $O(N^4)$

- 2'825 atoms
  - 31'247 basis functions
  - 976'375'009 elements in P
  - 953'308'158'199'750'081 integrals
- Exa-Pet-Ter-Gig-Meg-Kil



$$O(N^4) \rightarrow O(N)$$

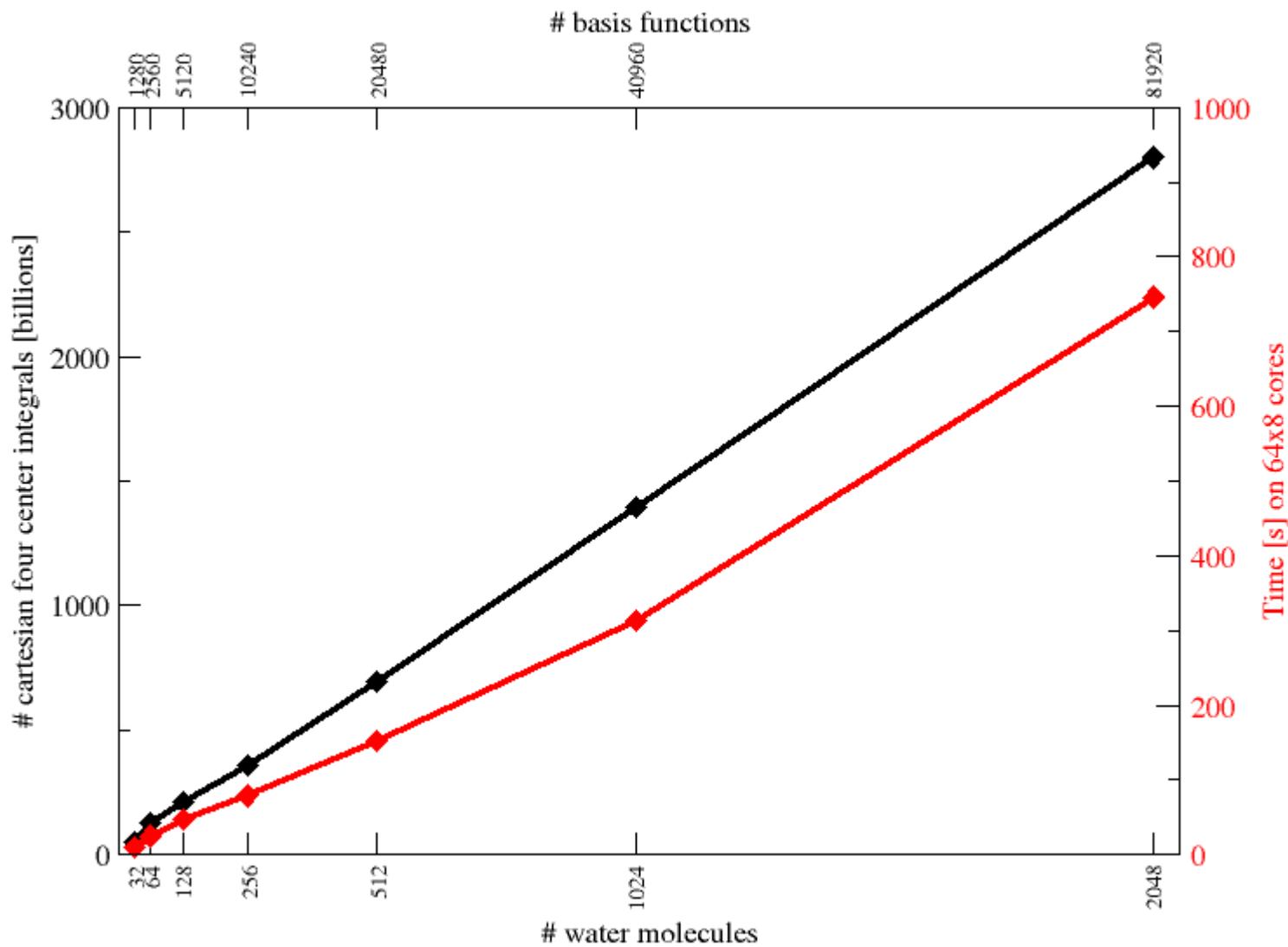
$$E_x^{\text{HF}} = -\frac{1}{2} \sum_{\alpha\beta\gamma\delta} P_{\alpha\beta} P_{\gamma\delta} (\phi_\alpha \phi_\gamma | \phi_\beta \phi_\delta)$$

$$(\phi_\alpha \phi_\gamma | \phi_\beta \phi_\delta) = \int d\mathbf{r} d\mathbf{r}' \frac{\phi_\alpha(\mathbf{r}) \phi_\gamma(\mathbf{r}) \phi_\beta(\mathbf{r}') \phi_\delta(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|}$$

Based on the fact that for large systems either the integrals  
Or the density matrix become zero (to within a threshold eps)

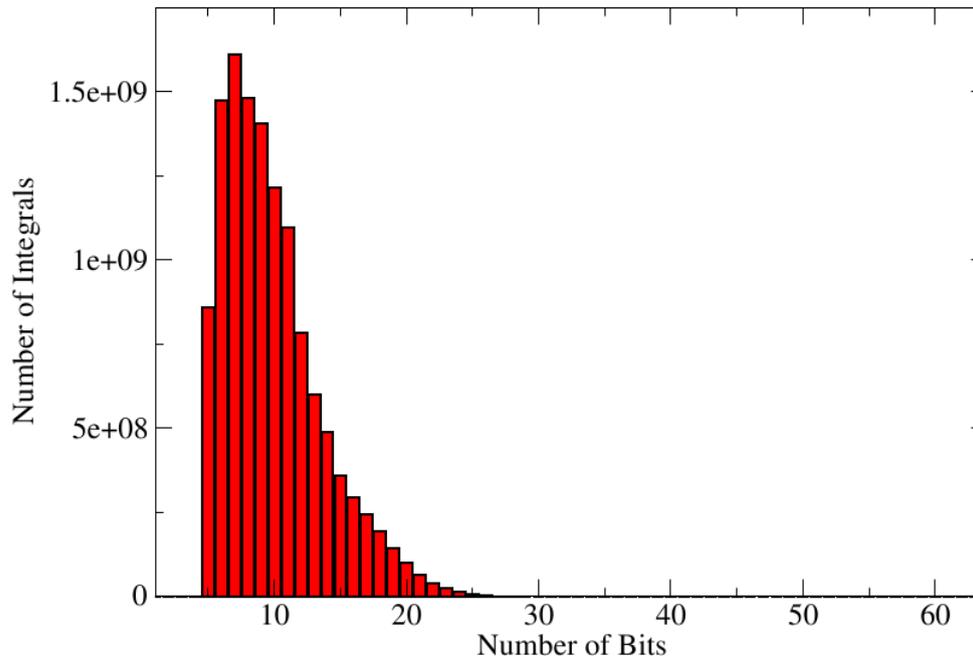
Cauchy-Schwarz screening	$ (ab cd)  \leq \sqrt{(ab ab)(cd cd)}$	$O(N^2)$
Density matrix screening	$P_{\alpha\beta}$ decays exponentially	$O(N)$
Operator screening	Operators other than $1/r$	$O(N)$

# $O(N)$ HFX: measurements



Linear scaling is key .... thousands of molecules possible  
On 'standard' cluster hardware in minutes.

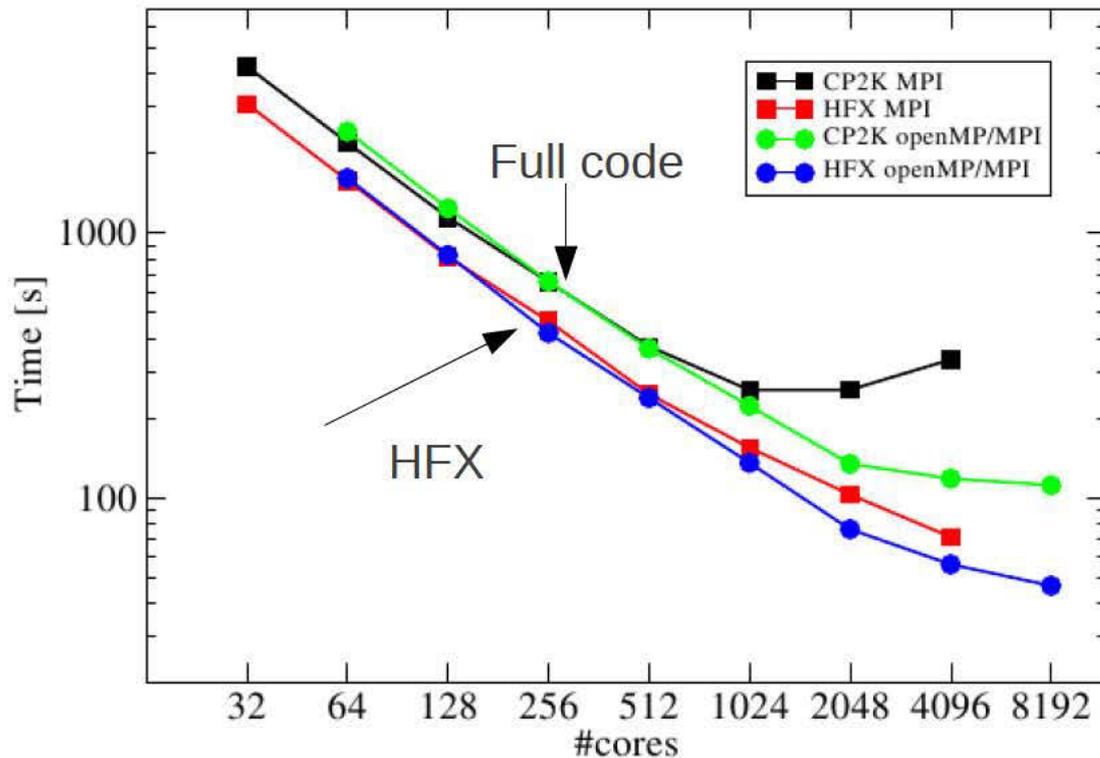
# In-core integral compression



Almost all simulations are performed using an in-core algorithm  
→ 10x speedup is observed.  
Highly efficient scheme: index free and lossy compression

# Parallel efficiency

HFX remains computationally much more demanding than GGA DFT (10x?)  
A good parallel implementation is mandatory



10 steps of MD, 64 H<sub>2</sub>O, 2560 BF,  
OpenMP: 8 threads/node

HFX code out-scales the  
Remaining (GGA) part of CP2K

Provided enough compute power,  
Hybrid simulations run essentially  
as fast as GGA  
(9s / BOMD step @ 4096 cores)

# Robust Hartree-Fock exchange in the condensed phase

$$E_x^{\text{PBC}} = -\frac{1}{2N_k} \sum_{i,j} \sum_{\mathbf{k},\mathbf{k}'} \int \int \psi_i^{\mathbf{k}}(r_1) \psi_j^{\mathbf{k}'}(r_1) g(|r_1 - r_2|) \psi_i^{\mathbf{k}}(r_2) \psi_j^{\mathbf{k}'}(r_2) d^3 r_1 d^3 r_2$$

How to treat this expression,  $\mathbf{k}=\mathbf{k}'$  is only conditionally convergent for  $g(r)=1/r$  ?

This 'difficult' point is integrable, but what for a Gamma-point code ( $\mathbf{k}=\mathbf{k}'=0$ ) ?

$$g_{\text{TC}}(r_{12}) = \begin{cases} \frac{1}{r_{12}}, & r_{12} \leq R_c \\ 0, & r_{12} > R_c \end{cases}$$

Avoids spurious self-exchange interactions with images in other cells.

This can be implemented robustly in a simple way.

Truncated Coulomb: Needs, Alavi

# LiH: demonstrating robustness & accuracy

A real benchmark system... a challenge inspired by the success of Quantum Chemistry

what is the final answer (i.e. many digits) for the HFX cohesive energy of LiH ?

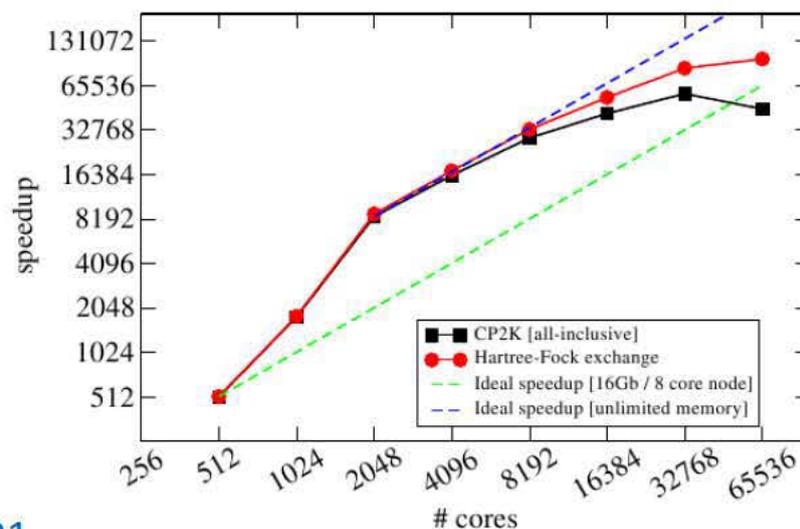
High quality basis (39000 functions for  $5 \times 5 \times 5 = 1000$  atoms)

	$R_c [\text{\AA}]$	$E(\text{HF})[\text{a.u.}]$	$H[\text{a.u.}](a)$	$\text{Li}[\text{a.u.}](b)$	$\epsilon_{\text{HF}}^{\text{coh}}[\text{a.u.}]$
2x2x2	4.0	-32.244609	-0.499957	-7.428493	-0.132702
3x3x3	6.0	-32.256844	-0.499974	-7.432137	-0.132100
4x4x4	8.0	-32.258022	-0.499974	-7.432582	-0.131949
5x5x5	10.0	-32.258179	N/A	N/A	N/A

vs. -0.13195 by Gillan, Manby, et al

vs. -32.258171 by Scuseria et al

Highly accurate calculations at the basis set limit, at gamma are possible... also for systems containing 1000 atoms



Paier J; Diaconu CV; Scuseria GE; Guidon M; VandeVondele J; Hutter J. 2009: PRB 80(17): 174114

Guidon M; Hutter J; VandeVondele J. 2009: JCTC 5(11): 3010-3021

# The curse of HFX: The Basis

Good quality calculations need good quality basis sets



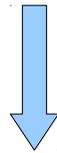
Increasing the quality of the basis steeply [ $O(N^4)$ ] increases the cost of HFX

1<sup>st</sup> HFX step for 20 water molecules on 128 cores

basis	$\kappa(S)$	threshold	cost [ERIs]	cost [s]	
3-21G*	4.9E+01	1.0E-04	2.3E+07	0.06	
6-31G**	2.1E+02	1.0E-05	5.2E+08	0.35	→ Most commonly used basis
6-311G++G**	1.2E+05	1.0E-07	1.1E+10	11.71	→
pc-0	5.2E+01	1.0E-04	1.7E+07	0.07	→
pc-1	4.5E+03	1.0E-05	4.4E+08	0.50	→ 'good' results
pc-2	5.7E+05	1.0E-07	2.0E+10	11.21	→
aug-pc-1	1.4E+06	1.0E-08	5.0E+10	53.23	→
aug-pc-2	3.9E+08	1.0E-09	1.5E+12	766.92	→
def2-QZVP	7.1E+04	1.0E-08	3.2E+11	127.16	→
aug-def2-QZVP	8.5E+05	1.0E-08	6.2E+11	331.61	→ 'Converged' results

# Auxiliary Density Matrix Methods (ADMM)

For certain density matrices HFX can be computed very efficiently  
(e.g. small basis sets or increased sparsity)



Transform an expensive matrix into a cheap one,  
use a GGA for estimating the difference

$$\begin{aligned} E_x^{\text{HFX}}[P] &= E_x^{\text{HFX}}[\hat{P}] + (E_x^{\text{HFX}}[P] - E_x^{\text{HFX}}[\hat{P}]) \\ &\approx E_x^{\text{HFX}}[\hat{P}] + (E_x^{\text{DFT}}[P] - E_x^{\text{DFT}}[\hat{P}]) \end{aligned}$$

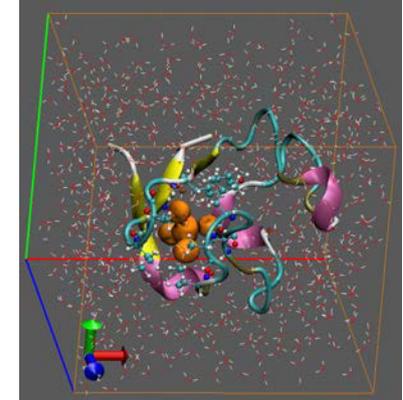
One example: wavefunction fitting, using an auxiliary basis

$$\min_{\tilde{C}} \left[ \sum_j \int (\psi_j(\mathbf{r}) - \tilde{\psi}_j(\mathbf{r}))^2 d\mathbf{r} + \sum_{k,l} \Lambda_{kl} \left( \int \tilde{\psi}_k(\mathbf{r}) \tilde{\psi}_l(\mathbf{r}) d\mathbf{r} - \delta_{kl} \right) \right]$$

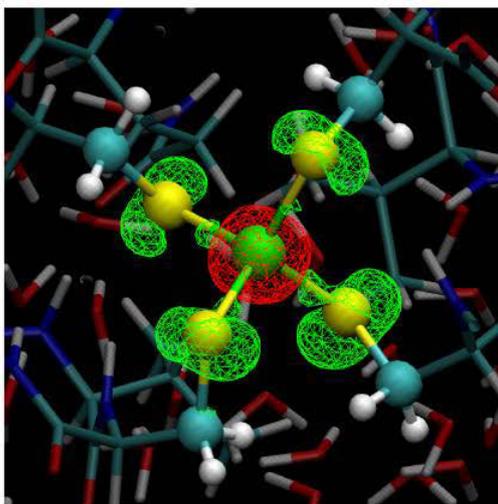
# ADMM: accuracy

- GTMKN24 (>1000 datapoints)
  - using FIT3 and pFIT3 auxiliary basis sets
    - Error wrt to Exp. : identical (5.0 kcal/mol)
    - Error wrt to Ref. : < 1 kcal/mol
- BSSE tests
  - BSSE reduced by 5x compared to no ADMM
  - H<sub>2</sub>O dimer: 0.5 kcal/mol FIT3, 0.2 kcal/mol aug-FIT3

# ADMM: performance

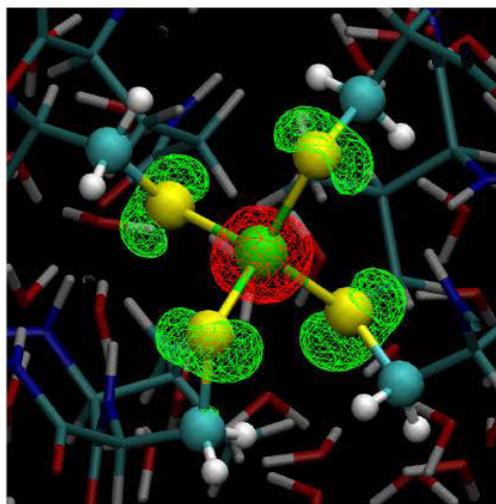


STD DZVP-MOLOPT-SR-GTH



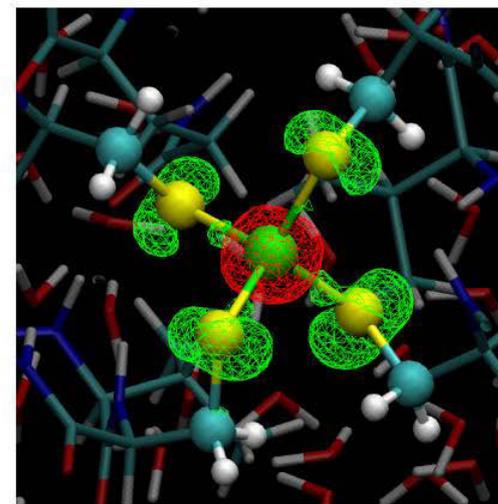
- ▶ 48'000 cores
- ▶ 1. SCF step: 45 min
- ▶ subsequent steps: 25 s
- ▶ 6.8 TB RAM

ADMM: MOLOPT/FIT3



- ▶ 1'152 cores
- ▶ 1. SCF step: 75 s
- ▶ subsequent steps: 25 s
- ▶ 5.2 GB RAM

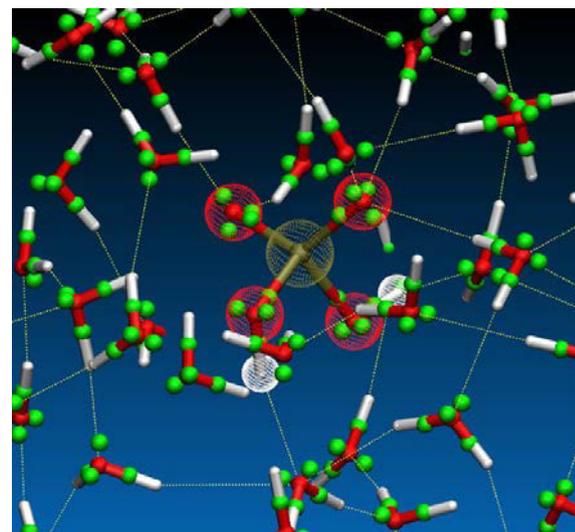
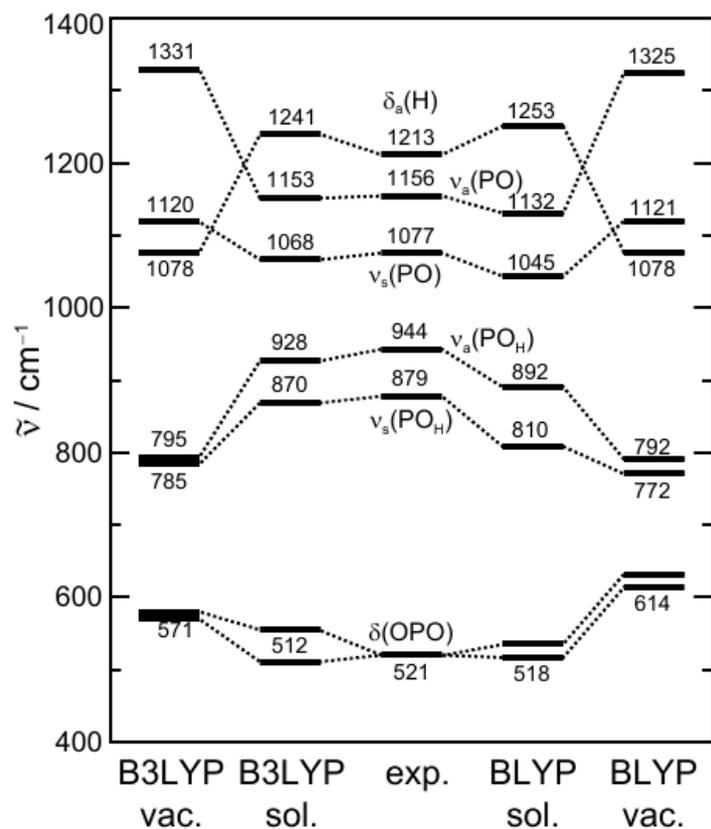
ADMM: MOLOPT/EMBED



- ▶ 1'152 cores
- ▶ 1. SCF step: 80 s
- ▶ subsequent steps: 25 s
- ▶ 5.2 GB RAM

A fully solvated protein computed within minutes using hybrid functionals

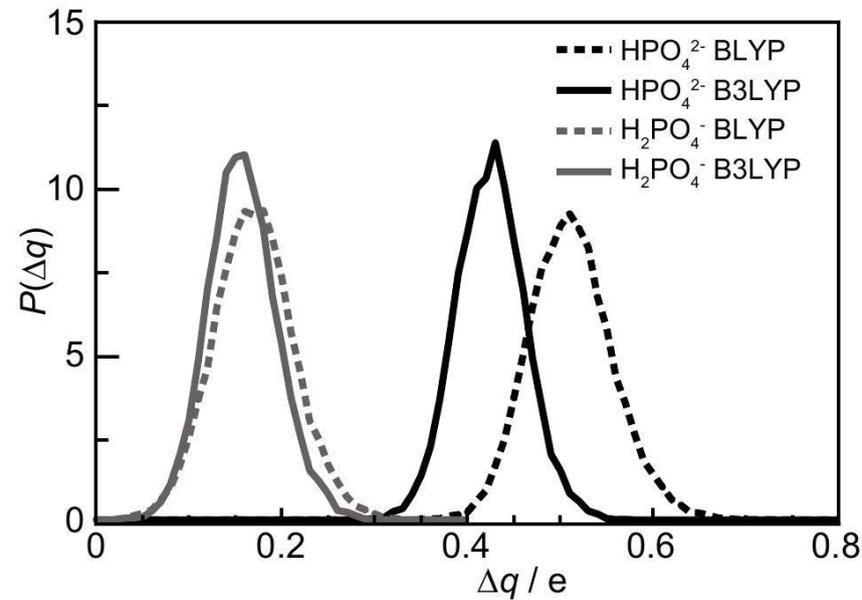
# IR spectroscopy from AIMD with hybrid functionals



~ 100 ps AIMD / 64 waters

Both Hybrids and GGA capture the main effect of solvation.  
Hybrid give 2x more accurate spectra

# Phosphate in solution



Analyzing the amount of charge transfer in  $\text{H}_2\text{PO}_4^-$  and  $\text{HPO}_4^{2-}$ :  
Reduced charge transfer is observed with hybrid functionals,  
but only for the double anion.

The effect is partially electronic, partially geometric

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Florian Schiffmann  
Atsushi Urakawa  
Ronny Wirz  
Alfons Baiker

## Cambridge

Michiel Sprik  
Marialore Sulpizi  
Jun Cheng

## PNL&Minnesota

Chris Mundy  
Roger Rousseau  
Ilja Siepmann

## World-wide

CP2K Team

## Flops&More

CSCS  
UZH  
SNF  
INCITE  
EU-FP  
DEISA  
PRACE

## UCL&EPCC&CRAY

Iain Bethune  
Matt Wattkins  
Ben Slater  
John Levesque

## München

Paul Tavan  
Gerald Matthias

**You for your attention!**